

MSSM4G[📶] scenario and vectorlike lepton @ LHC

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Based on [[1608.00283](#)] in collaboration with
M. Abdullah, J. L. Feng, and B. Lillard (UC Irvine)

chiral: $\mathcal{L} \ni y \langle H \rangle \bar{\psi}_R \psi_L$ vector-like: $\mathcal{L} \ni m \bar{\psi}_R \psi_L$

■ Chiral q_4 ... disfavored by
 $\sigma(gg \rightarrow h)$

■ Vectorlike q_4

- Motivated by MSSM Higgs mass enhancement
- Searched for at the LHC: $pp \rightarrow q_4 \bar{q}_4 \rightarrow 2q + 2 \text{ bosons}$

■ Chiral e_4 ... ??? (at least GUT incompatible?)

■ Vectorlike e_4

- Motivated by MSSM4G model (and more)
- Search-able at the (HL-)LHC

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- Search-able at the (HL-)LHC **3**

■ SUSY = solution to Naturalness

old-days expectation :

Papucci, Ruderman, Weiler [[1110.6926](#)]

$$|\mu| \lesssim 200 \text{ GeV} \left(\frac{\Delta^{-1}}{20\%} \right)^{-1/2}$$

$$\sqrt{m_{\tilde{t}_1}^2 + m_{\tilde{t}_2}^2} \lesssim 600 \text{ GeV} \frac{\sin \beta}{\sqrt{1 + \alpha^2}} \left(\frac{\log(\Lambda/\text{TeV})}{3} \right)^{-1/2} \left(\frac{\Delta^{-1}}{20\%} \right)^{-1/2}$$

$$m_{\tilde{g}} \lesssim 900 \text{ GeV} \cdot \sin \beta \left(\frac{\log(\Lambda/\text{TeV})}{3} \right)^{-1/2} \left(\frac{\Delta^{-1}}{20\%} \right)^{-1/2}$$

difficulties:

- LHC bounds
- radiative Higgs mass

Δ : “naturalness”

$\alpha = A_t/m_{\tilde{t}}$ (stop mixing parameter)

Λ : SUSY-breaking scale

■ SUSY = solution

old-days expectation :

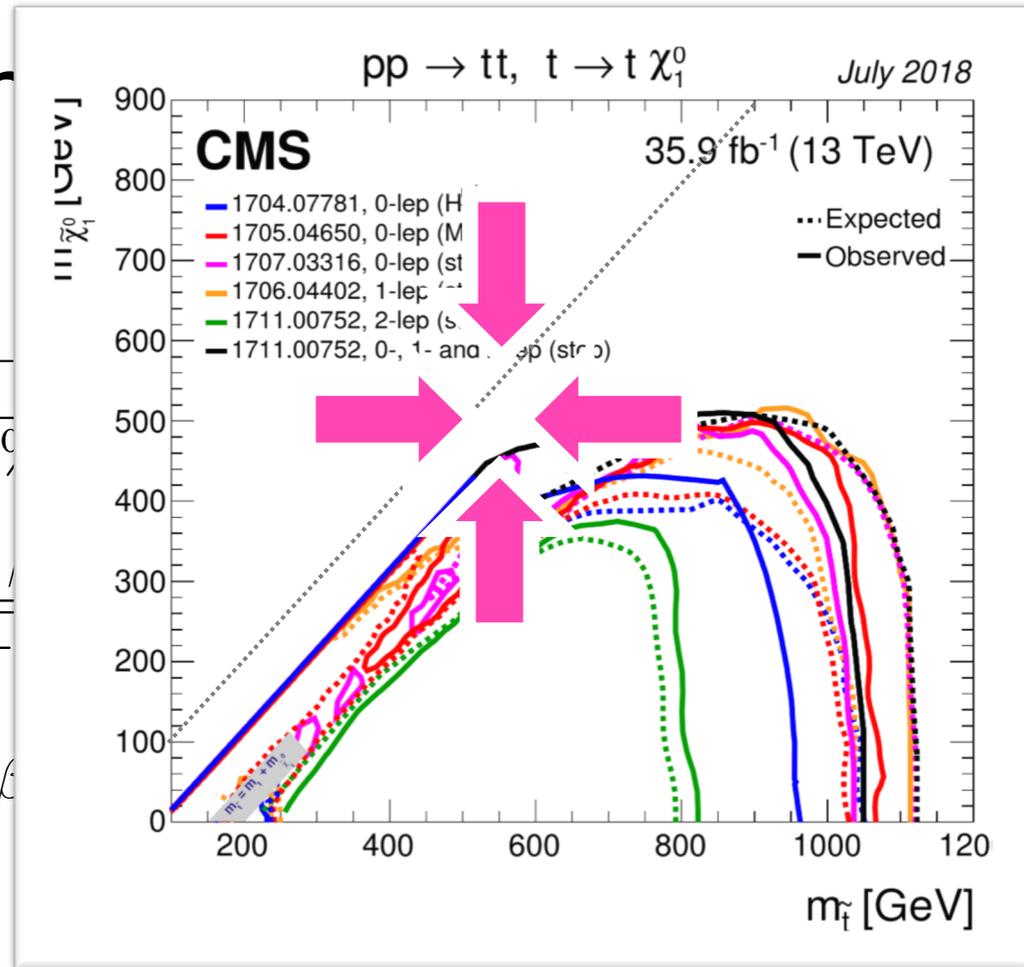
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difficulties:

- LHC bounds ... "avoidable"
- radiative Higgs mass



Δ : "naturalness"

$a = A_t/m_{\tilde{t}}$ (stop mixing parameter)

Λ : SUSY-breaking scale

■ MSSM prediction

$$m_h^2 \approx m_Z^2 \cos^2(2\beta) + \frac{3y_t^2 \sin^2 \beta}{4\pi^2} m_t^2 \left[\log \frac{m_{\tilde{t}}^2}{m_t^2} + \alpha^2 \left(1 - \frac{\alpha^2}{12} \right) \right]$$

$(125\text{GeV})^2$ $(91\text{GeV})^2$ SUSY-breaking effect from $W \ni \underline{y_t H_u q_3 \bar{u}_3}$

$$\implies m_{\tilde{t}} \gtrsim 1 \text{ TeV}$$

cf. Hall, Pinner, Ruderman [[1112.2703](#)] etc.

■ To relax this constraint

- add a scalar particle ("NMSSM")
- add "another top" ... MSSM + vectorlike quark

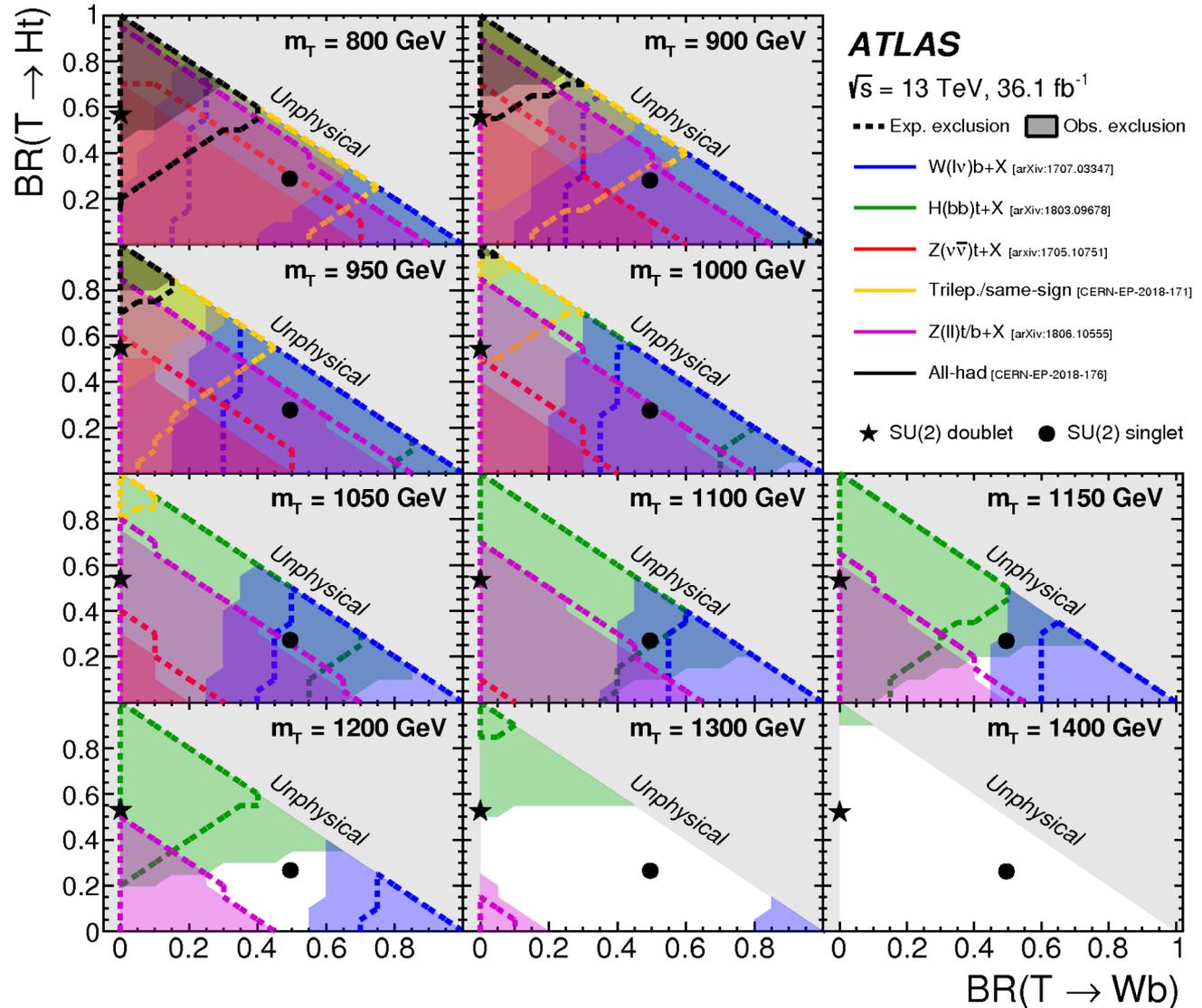
e.g. MSSM + $q_4 + \bar{u}_4 + \bar{q}_4 + u_4$:

$$W_{\text{VLQ}} = \underline{y' H_u q_4 \bar{u}_4} + \underline{M_{q_4} q_4 \bar{q}_4} + \underline{M_{u_4} u_4 \bar{u}_4}$$

vectorlike mass terms

assuming $T \rightarrow Wb$ or Ht or Zt ,

$m_T < 1100$ GeV is excluded for any decay patterns (BRs).



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e.g. MSSM + $q_4 + \bar{u}_4 + \bar{q}_4 + u_4$: GUT compatibility?

$$W_{\text{VLQ}} = y' H_u q_4 \bar{u}_4 + M_{q_4} q_4$$

$$5 = q + \bar{u} + \bar{e} \text{ in SU(5)-GUT}$$

- MSSM4G scenario

= use this e_4 to solve "bino overabundance"

Abdullah, Feng [1510.06089]

If dark matter is $\tilde{\chi}^0$ and it is pure-Bino \tilde{B} ,
it would give larger DM density than observed.

$\tilde{\chi}^0$: lightest neutralino

$$\tilde{\chi}^0 = \tilde{B} \oplus \tilde{W}^0 \oplus \tilde{H}_d^0 \oplus \tilde{H}_u^0$$

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1) In early Universe with $T > m_{\tilde{B}}$ → DM "equilibrium"

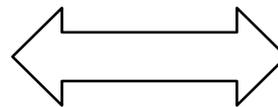
pair creation ($f\bar{f} \rightarrow \tilde{B}\tilde{B}$)



pair annihilation ($\tilde{B}\tilde{B} \rightarrow f\bar{f}$)



equilibrium



kinetic energy $\sim T$ ($> m_{\tilde{B}}$)

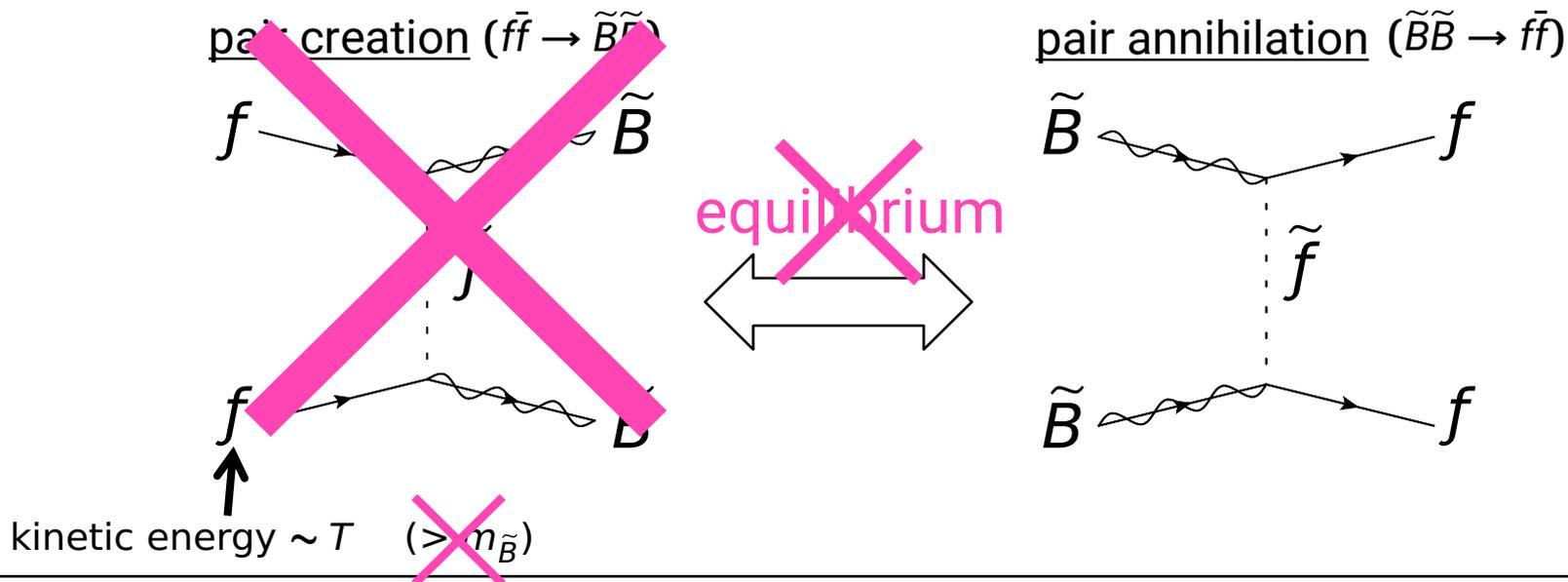
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2) In early Universe after $T < m_{\tilde{B}}$ → DM decreases



■ MSSM4G scenario

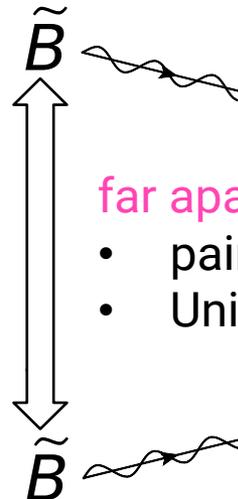
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3) after $T < m_{\tilde{B}}/20 \rightarrow$ DM **frozen-out**

pair annihilation ($\tilde{B}\tilde{B} \rightarrow f\bar{f}$)



- far apart due to
- pair annihilation
 - Universe's expansion

■ MSSM4G scenario

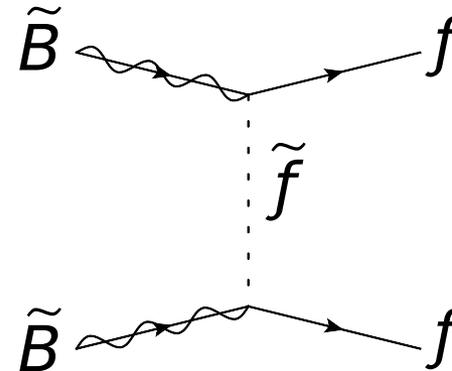
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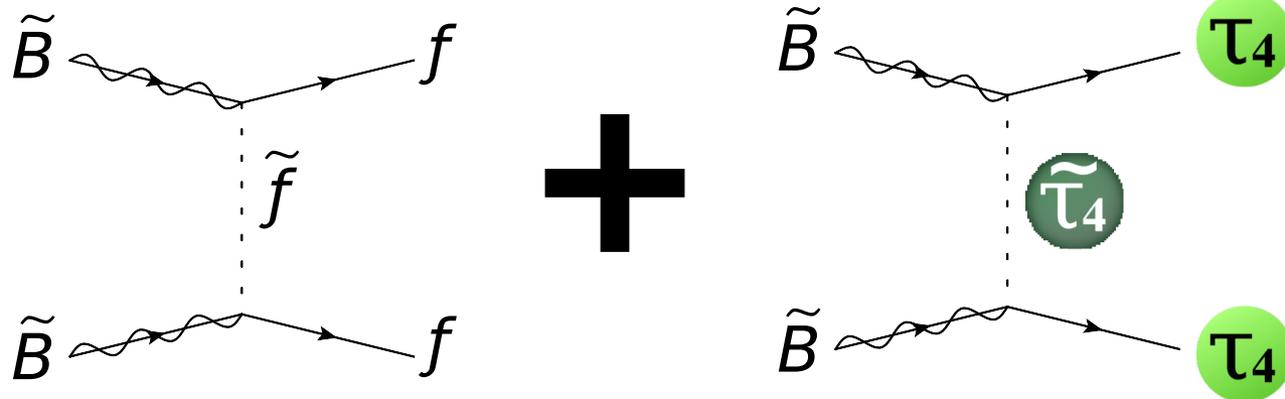


\tilde{B} : smaller cross section
→ "overabundant problem"
of Bino thermal relic DM

■ MSSM4G scenario

Abdullah, Feng [1510.06089]

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extra annihilation channel

if $\tau_4 \gtrsim \tilde{\tau}_4 > \tau_4$

→ abundance reduced

$$\langle \sigma v \rangle = \frac{g_Y^4 Y_L^2 Y_R^2}{2\pi} \frac{m_f^2}{m_{\tilde{B}}} \frac{\sqrt{m_{\tilde{B}}^2 - m_f^2}}{(m_{\tilde{B}}^2 + m_{\tilde{f}}^2 - m_f^2)^2}$$

■ MSSM4G scenario

= use this e_4 to solve "bino over

$Q : (\mathbf{3}, \mathbf{2}, 1/6)$	$\bar{Q} : (\bar{\mathbf{3}}, \mathbf{2}, -1/6)$
$\bar{U} : (\bar{\mathbf{3}}, \mathbf{1}, -2/3)$	$U : (\mathbf{3}, \mathbf{1}, 2/3)$
$\bar{D} : (\bar{\mathbf{3}}, \mathbf{1}, 1/3)$	$D : (\mathbf{3}, \mathbf{1}, -1/3)$
$\bar{E} : (\mathbf{1}, \mathbf{1}, 1)$ (same as SM)	$E : (\mathbf{1}, \mathbf{1}, -1)$ (vectorlike partners)

■ QUE model : MSSM + $Q\bar{Q}U\bar{U}E\bar{E}$

✓ gauge coupling unification

✓ SU(5) GUT

➤ extra $H_u Q_4 \bar{U}_4$ interaction $\rightarrow m_h$ 

■ QDEE model : MSSM + $Q\bar{Q}D\bar{D}E\bar{E}E\bar{E}$

✓ gauge coupling unification

~~✗ SU(5) GUT~~

➤ extra $H_d Q_4 \bar{D}_4$ coupling $\rightarrow m_h$ slightly 

■ MSSM4G scenario

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vectorlike mass terms

$$W_{\text{QUE}} = y' H_u Q_4 \bar{U}_4 - y'' H_d \bar{Q}_4 U_4 + \overbrace{M_Q Q_4 \bar{Q}_4 + M_U U_4 \bar{U}_4 + M_E E_4 \bar{E}_4}^{\text{vectorlike mass terms}}$$

$$+ \underbrace{\epsilon_i H_u Q_i \bar{U}_4 + \epsilon'_i H_u Q_4 \bar{U}_i - \epsilon''_i H_d L_i \bar{E}_4}_{\text{SM-Vectorlike mixings}}$$

SM-Vectorlike mixings

- **dangerous** : induce flavor violations
- **necessary** : let VLF decay

... we need an approximate

DISCRETE Z_2

(SM: even, VLF: odd)

CASAM
INVESTICAO



DISCRETE18

■ MSSM4G scenario

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$$W_{QDEE} = y' H_u \bar{Q}_4 D_4 - y'' H_d Q_4 \bar{D}_4 + M_Q Q_4 \bar{Q}_4 + M_D D_4 \bar{D}_4 + M_E E_4 \bar{E}_4 + M'_E E_5 \bar{E}_5$$

$$- \epsilon_i H_d Q_i \bar{D}_4 - \epsilon'_i H_d Q_4 \bar{D}_i - \epsilon''_i H_d L_i \bar{E}_4 - \epsilon'''_i H_d L_i \bar{E}_5$$

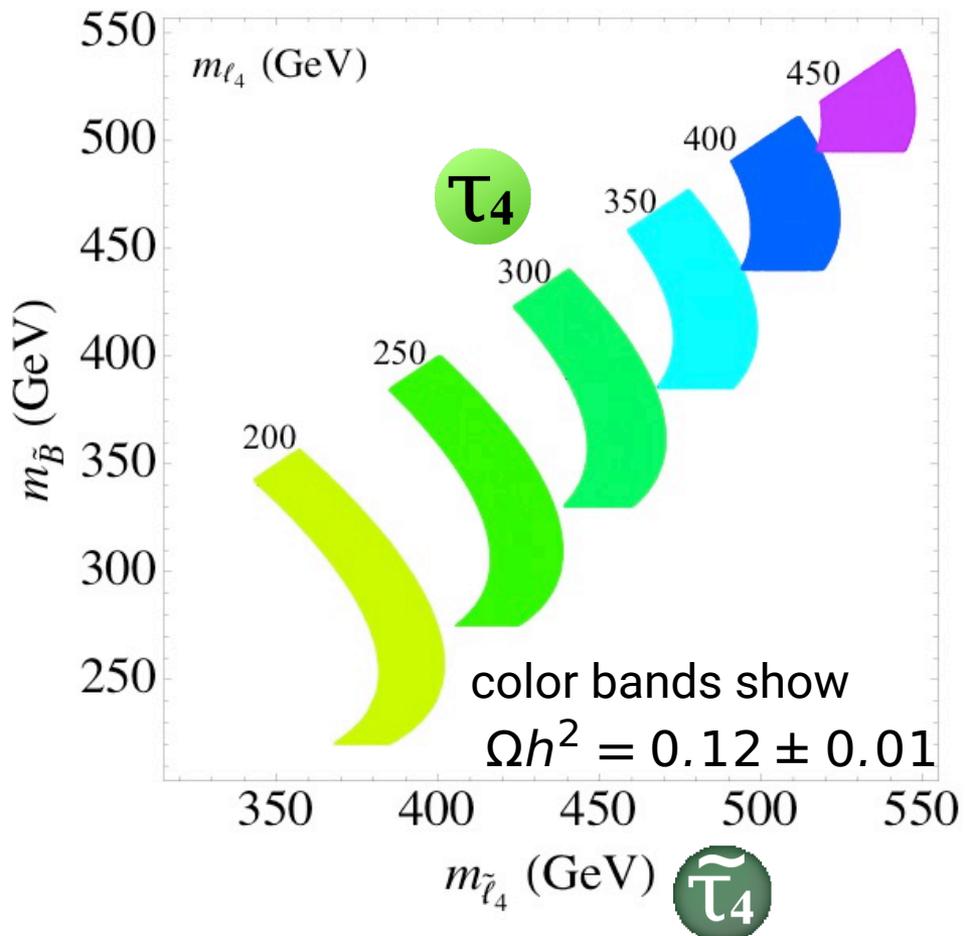
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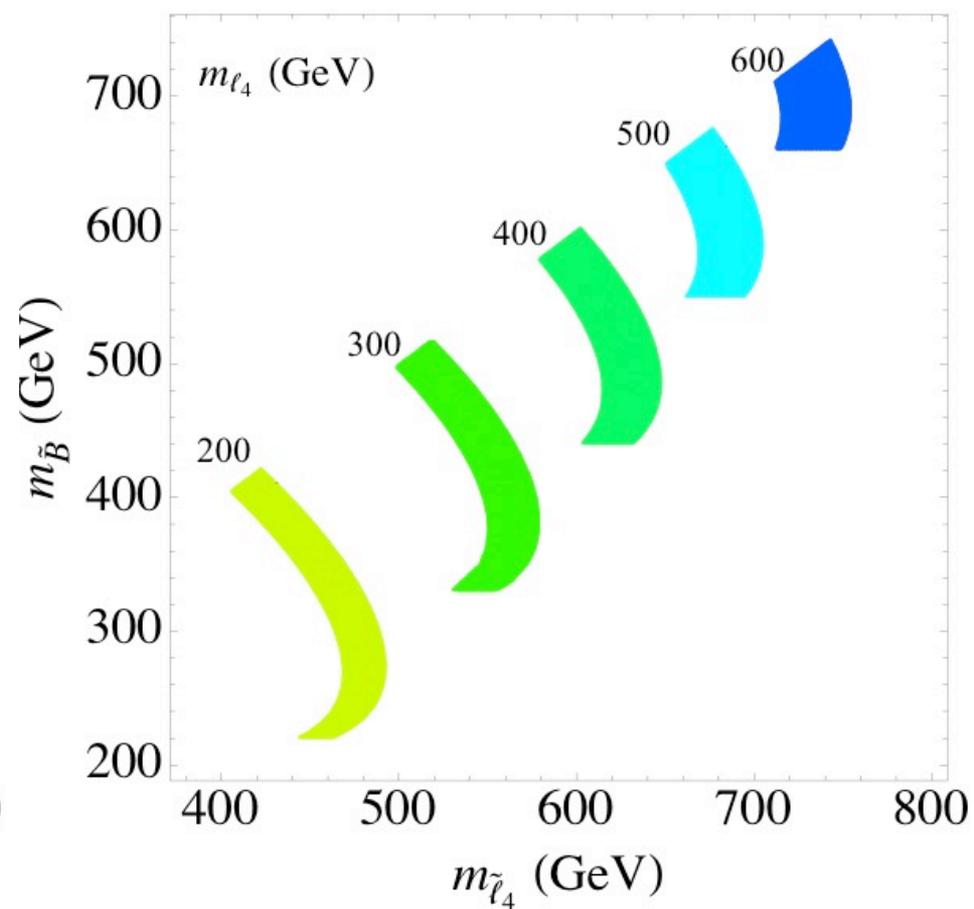
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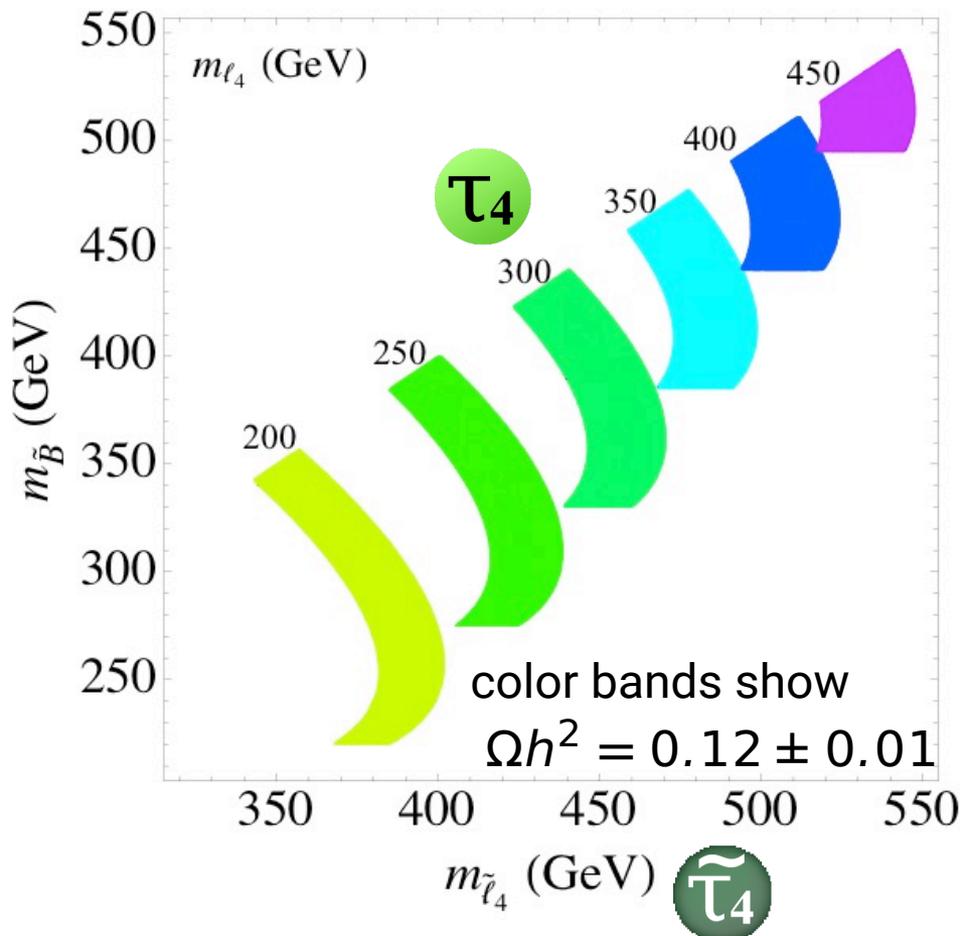


QDEE model



$$\tilde{\tau}_4 \gtrsim \tilde{B} > \tau_4$$

QUE model



■ LHC searches

- Searches for VL-slepton $\tilde{\tau}_4$
 - 2 lepton + mET search
- Searches for VLL τ_4

■ Gamma-ray observation

- from τ -leptons coming from
 - (DM)+(DM) → $\tau_4 \tilde{\tau}_4$
 - $(Z\tau)(Z\bar{\tau})$

■ (DM direct detection)

- LZ, DARWIN, ...

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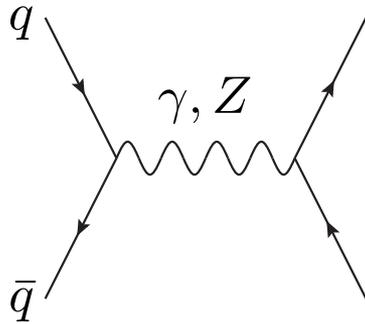
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■ Vector-like lepton τ_4

- Lagrangian $\mathcal{L}_{\text{extra}} = \epsilon_i H L_i \bar{\tau}_4 + m \bar{\tau}_4 \tau_4 + (\text{kinetic})$
- production & decay:

$$pp \rightarrow \tau_4 \bar{\tau}_4, \quad \tau_4 \rightarrow W^- \nu_i, \quad Z l_i^-, \quad h l_i^-$$

"Drell-Yan" production



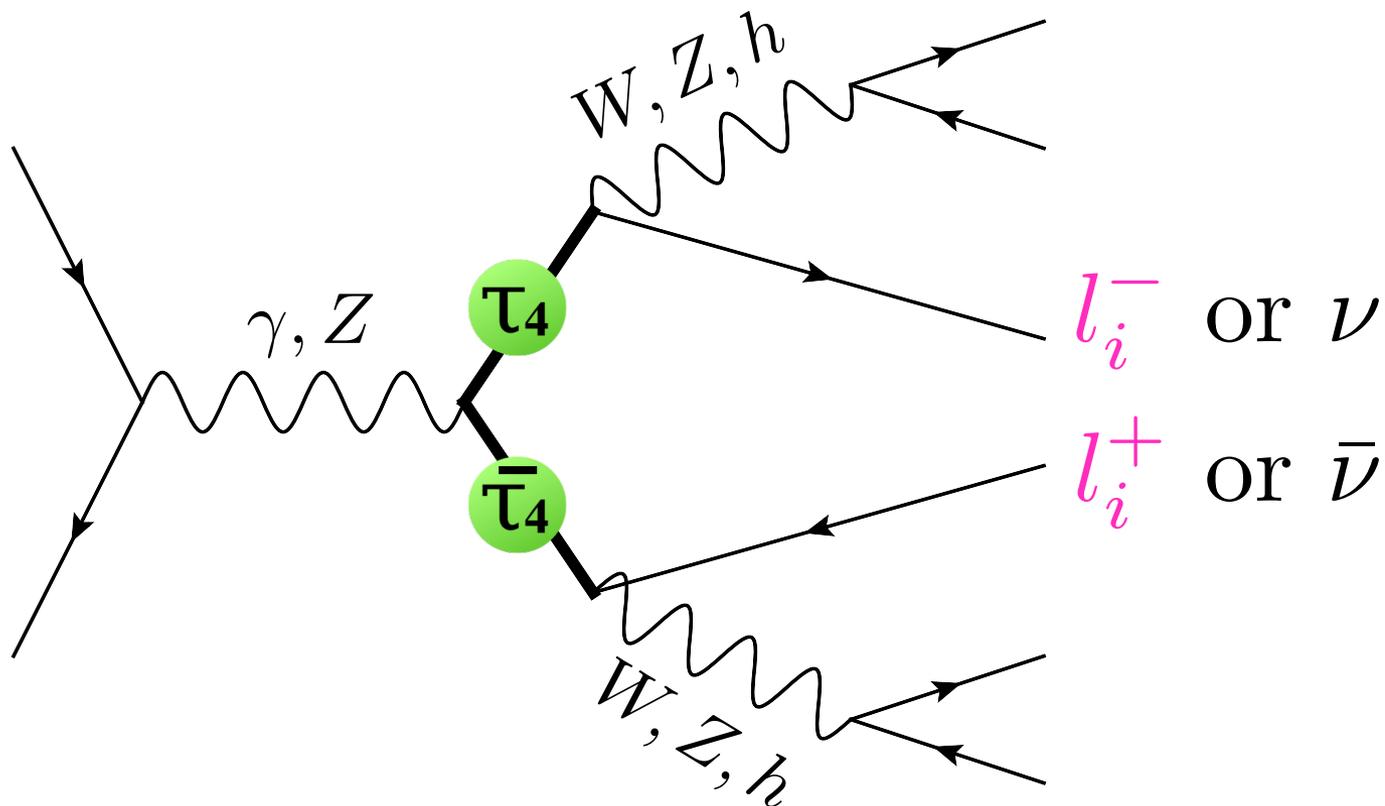
Decay through "SM-4G mixing"

- 3 benchmark scenarios assuming hierarchical ϵ_i
(if not \rightarrow flavor-violation constraints)

- $\text{Br}(\tau_4 \rightarrow e \text{ or } \nu_e) \approx 100\%$
- $\text{Br}(\tau_4 \rightarrow \mu \text{ or } \nu_\mu) \approx 100\%$
- $\text{Br}(\tau_4 \rightarrow \tau \text{ or } \nu_\tau) \approx 100\%$

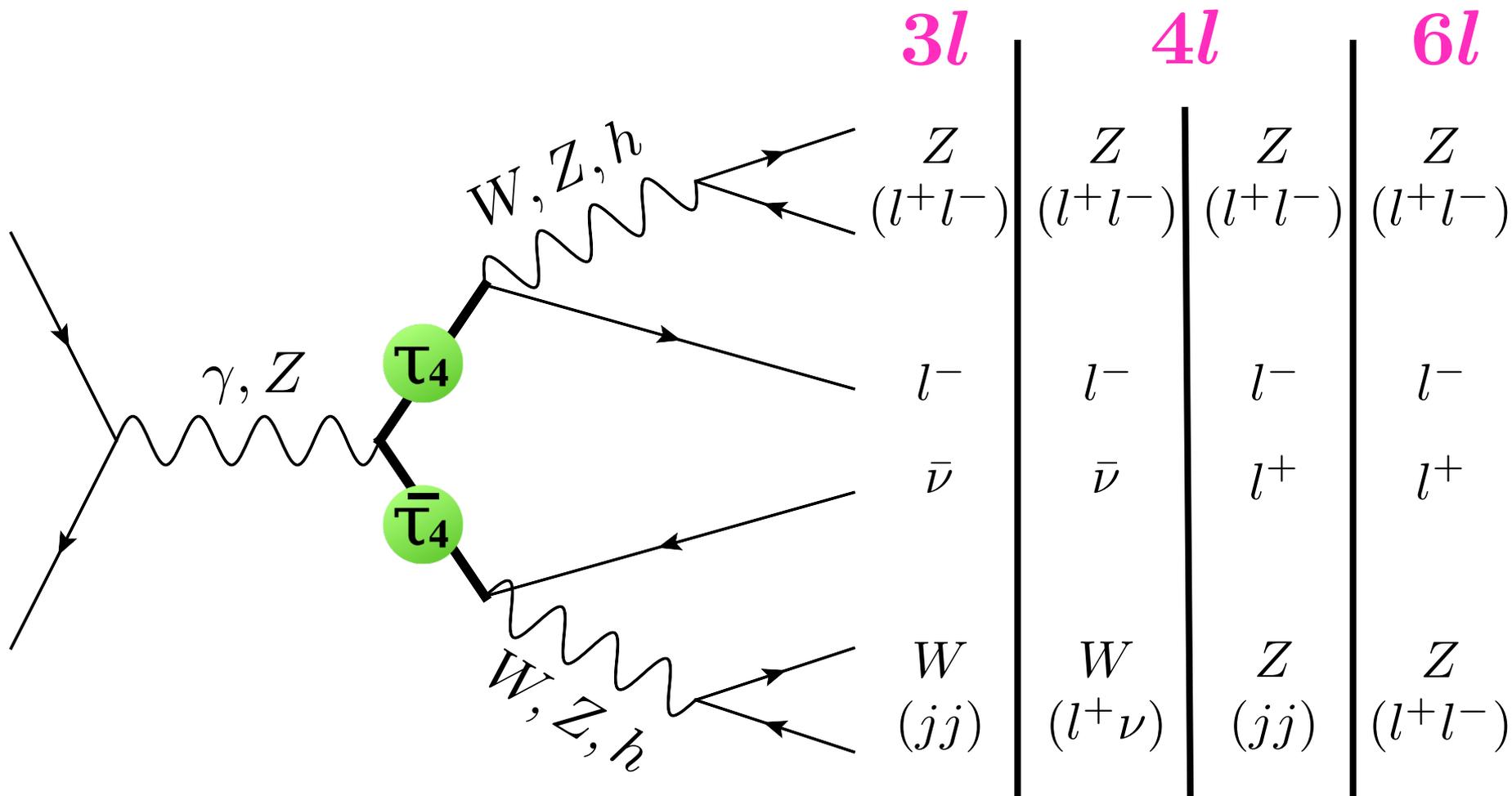
cf. $\text{Br}(W) : \text{Br}(Z) : \text{Br}(h) \approx 2 : 1 : 1$

- Vector-like lepton : $pp \rightarrow \tau_4 \bar{\tau}_4$ $\tau_4 \rightarrow Z l_i^-, W^- \nu_i, h l_i^-$



How do you capture this signature?

■ Vector-like lepton : $pp \rightarrow \tau_4 \bar{\tau}_4$ $\tau_4 \rightarrow Z l_i^-, W^- \nu_i, h l_i^-$



■ Vector-like lepton : $pp \rightarrow \tau_4 \bar{\tau}_4 \quad \tau_4 \longrightarrow Z l_i^-, W^- \nu_i, h l_i^-$

➤ Our "most basic" approach: **multi- ℓ^\pm (3–5 ℓ^\pm) search**

➔ HL-LHC may exclude $m_{\tau_4} < 350$ GeV
if VLL decays to e or μ .
 (= "mixes with")

details for experts

	$WZ(j)$	$WZ(\ell)$	$ZZ(j)$	$ZZ(\ell)$	
N_ℓ	≥ 3	≥ 4	≥ 4	≥ 5	
N_j	≥ 2	< 2	≥ 2	—	
$ m_{jj} - m_W $	< 20 GeV	—	—	—	
$ m_{jj} - m_Z $	—	—	< 40 GeV	—	←W/Z-like jet pair
\cancel{E}_T	> 60 GeV	> 100 GeV	—	—	←Large mET from ν &W
$N_{Z(\ell\ell)}$	—	—	≥ 1	≥ 1	←Z-like lepton pair

- Snowmass BKG set is used.
 - MG5–Pythia–Delphes + NLO K-factor
 - di-boson + tt dominated
- Signal by FR–MG5aMC–Pythia–Delphes (LO)

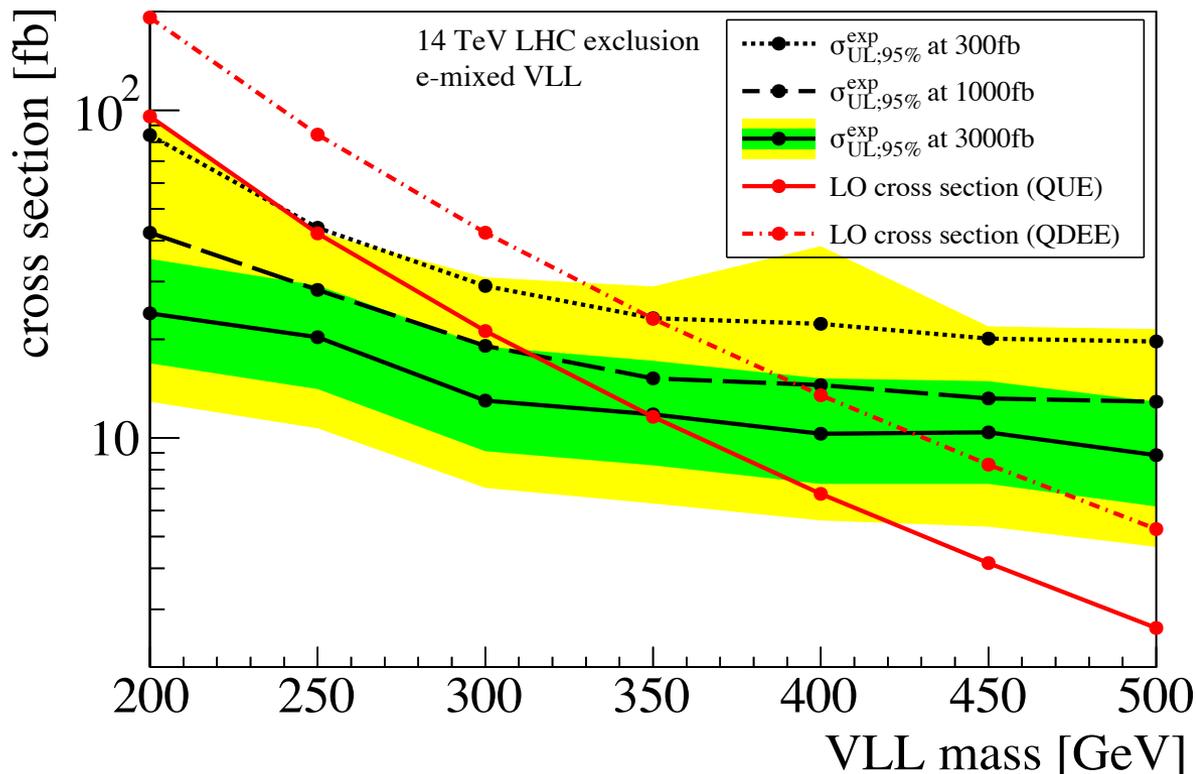
- PT cut: $(\ell_1, \ell_2, \ell_i) > (120, 60, 20)$ GeV, $(j) > 20$ GeV.
- **tau-tag / b-tag not used (avoided)**
- Uncertainties = stat. + 20% syst.

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■ Possible (by expm. physicists) improvements

➤ This bound is for "**right-handed**" VLL.

→ stronger bound expected for left-handed VLL.

➤ Use τ -tagging to capture VLL mixed with τ .

➤ Use Higgs-tagging to capture VLL decays to h .

■ Vectorlike q_4

- m_h (MSSM) ^{UP} ↗ ; excluded up to 1.1TeV

■ Vectorlike e_4

$$\mathcal{L}_{\text{extra}} = \epsilon_i H L_i \bar{\tau}_4 + m \bar{\tau}_4 \tau_4 + (\text{kinetic})$$

- MSSM4G model: solves Bino-overabundance
- HL-LHC expectation : **350 GeV** if "right-handed" & mixed with e or μ

What I skipped

- MSSM4G is testable also by **DM-direct** and **Gamma-ray**.

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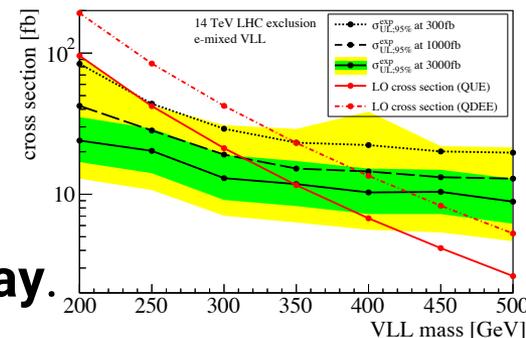
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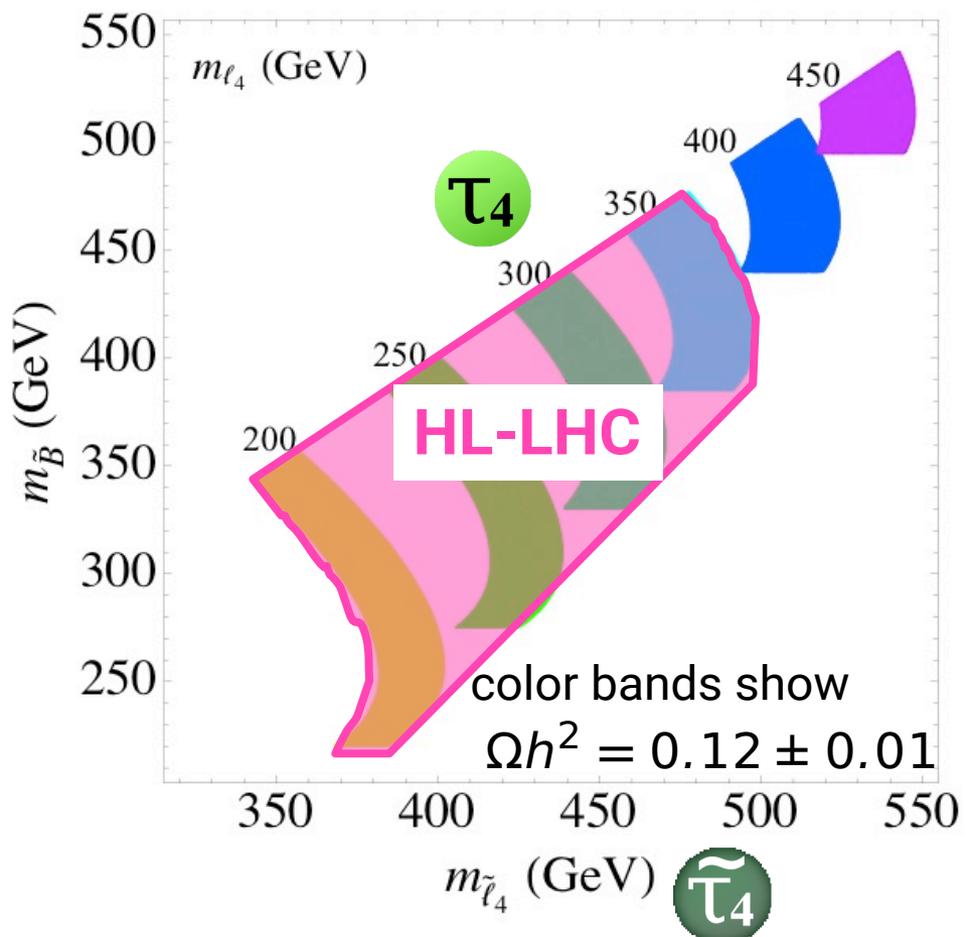
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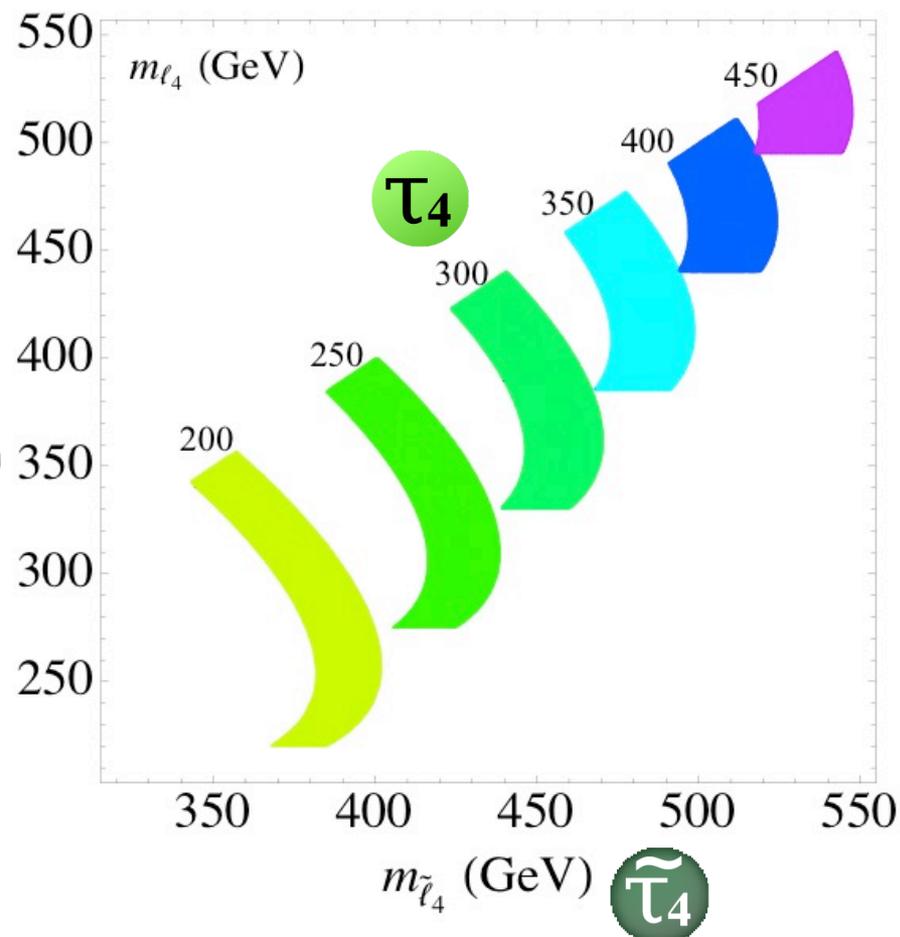
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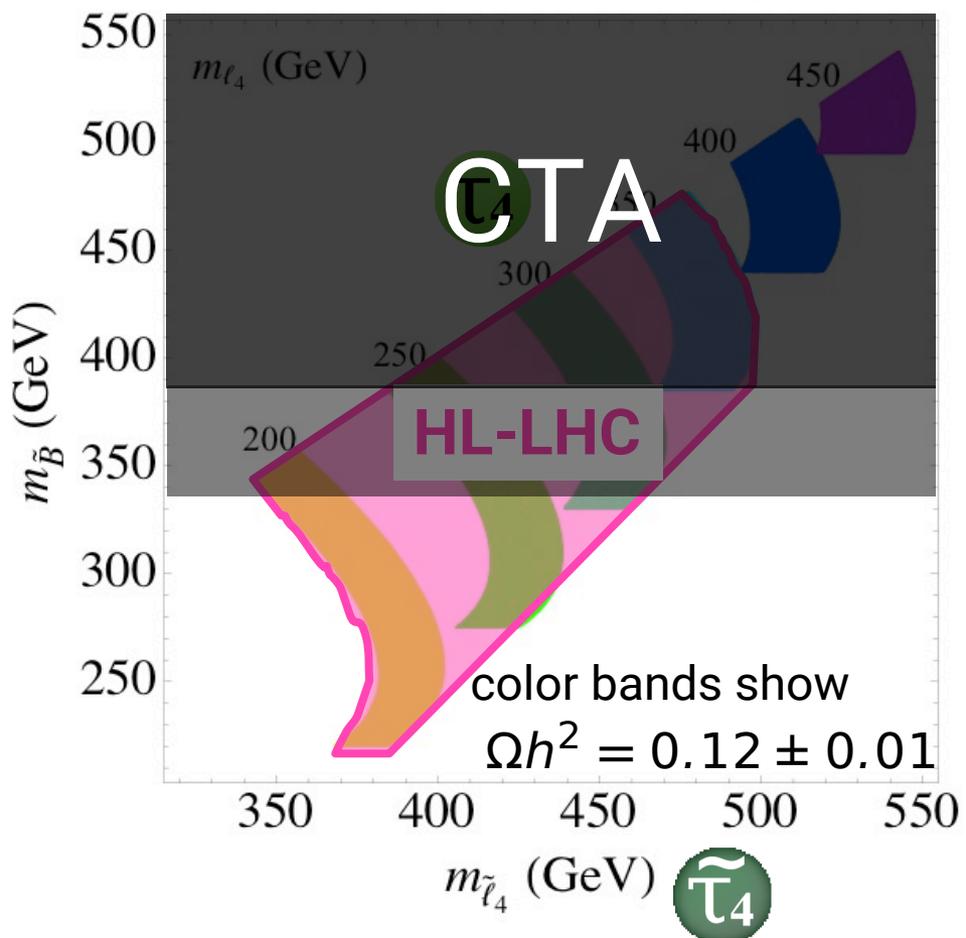
QUE model : e or μ



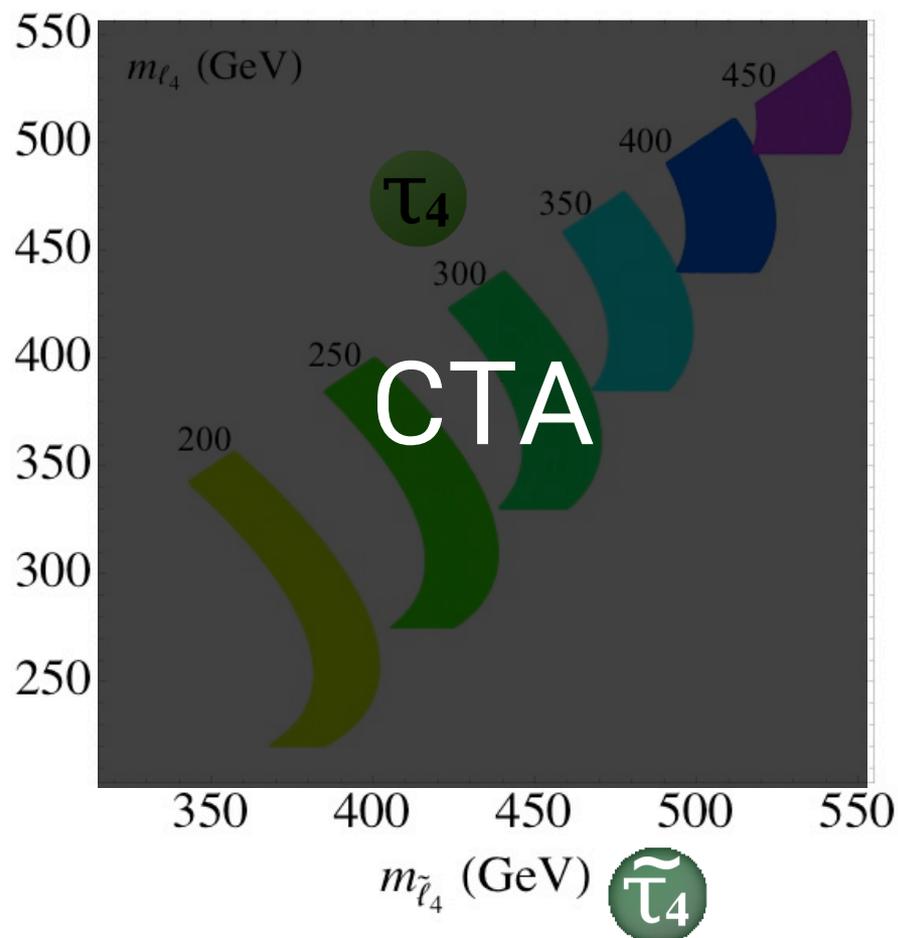
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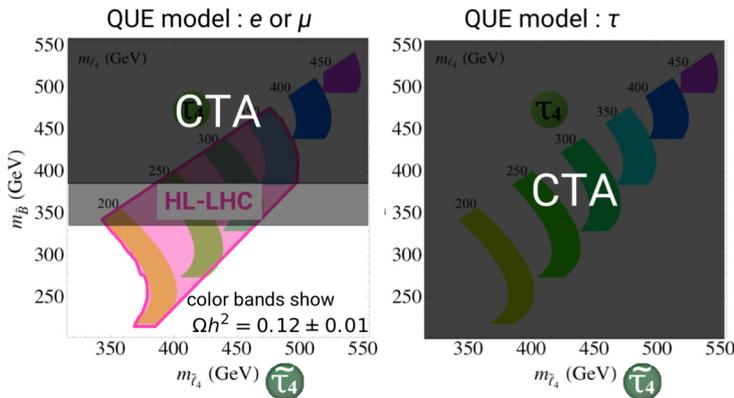
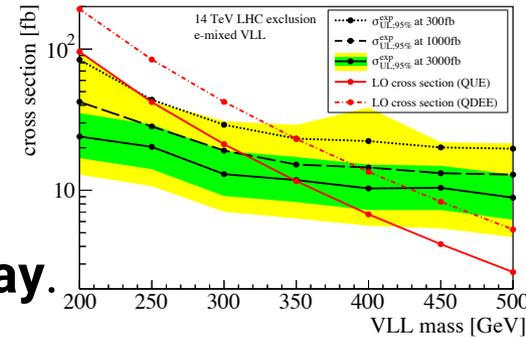
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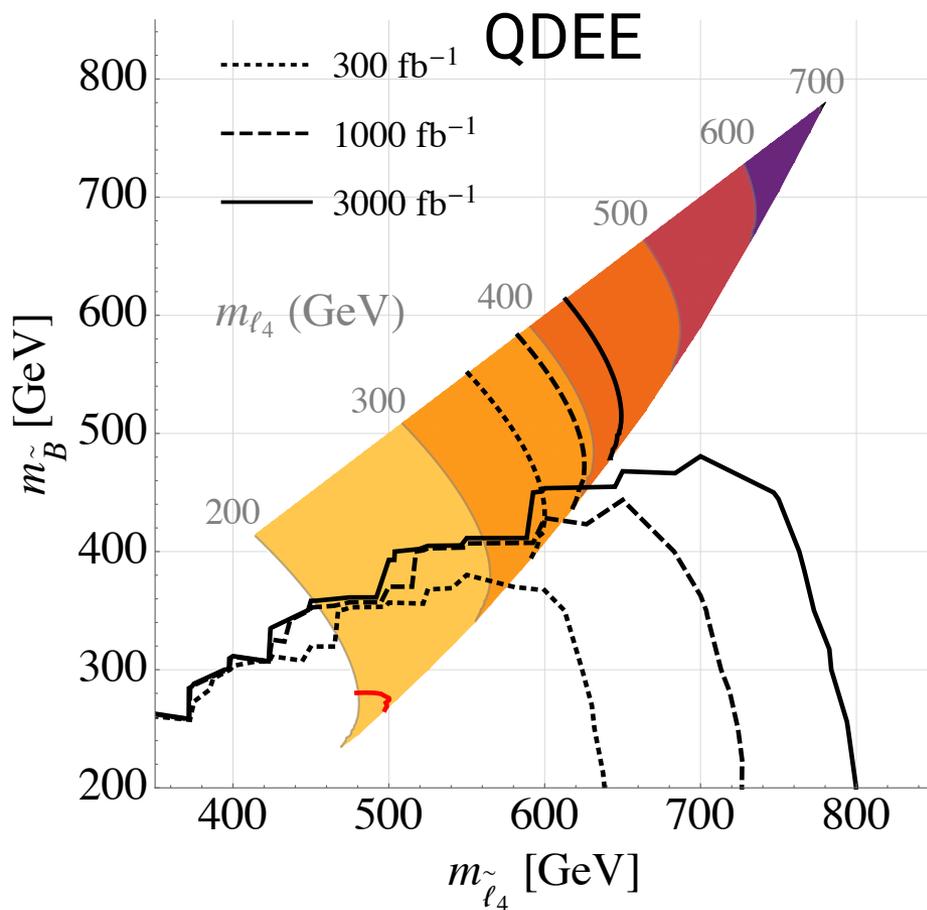
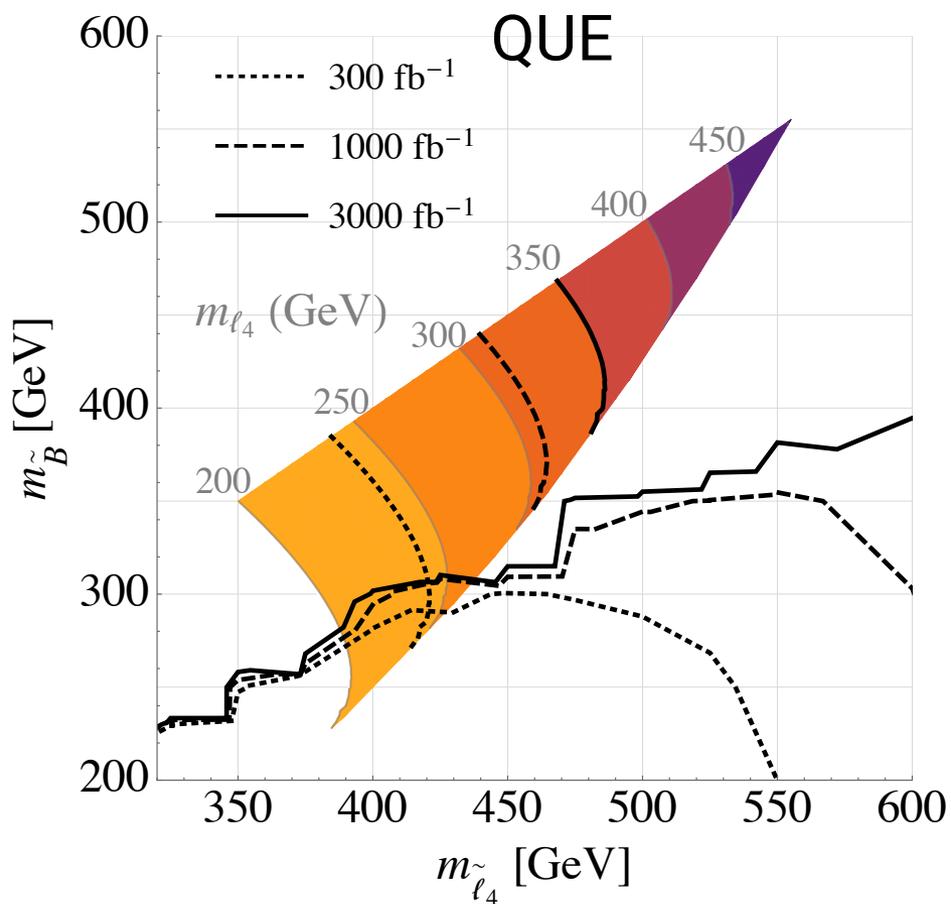
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MSSM4G "official" space

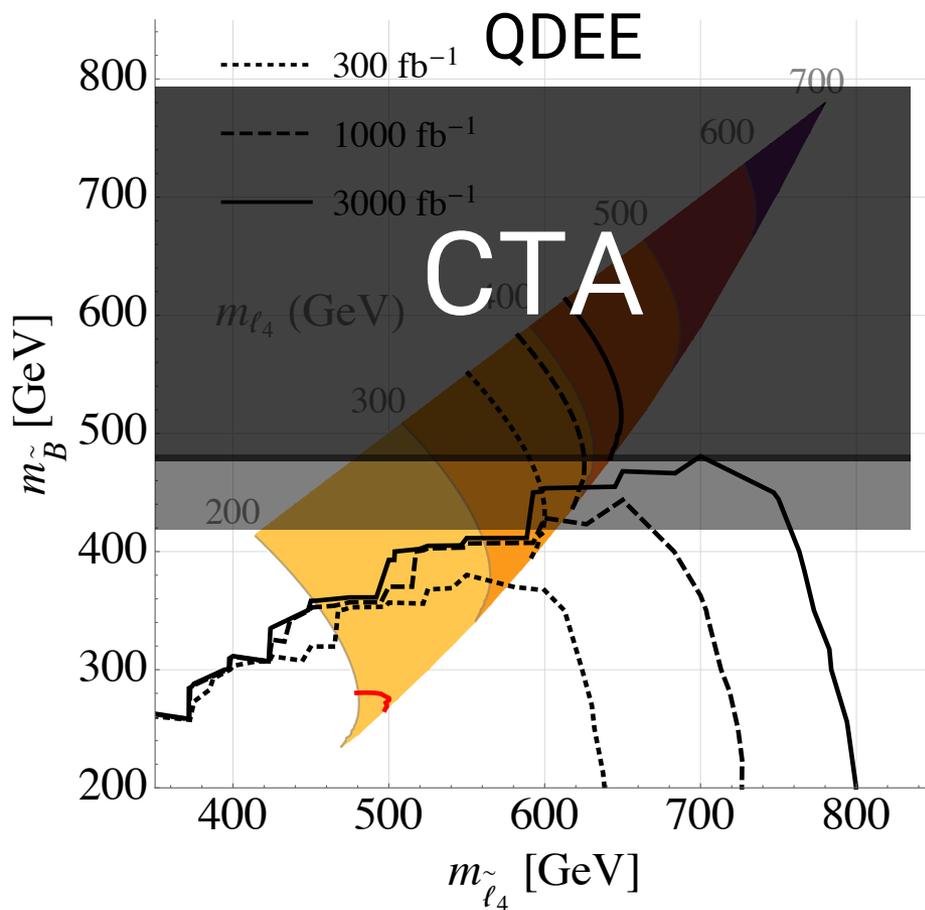
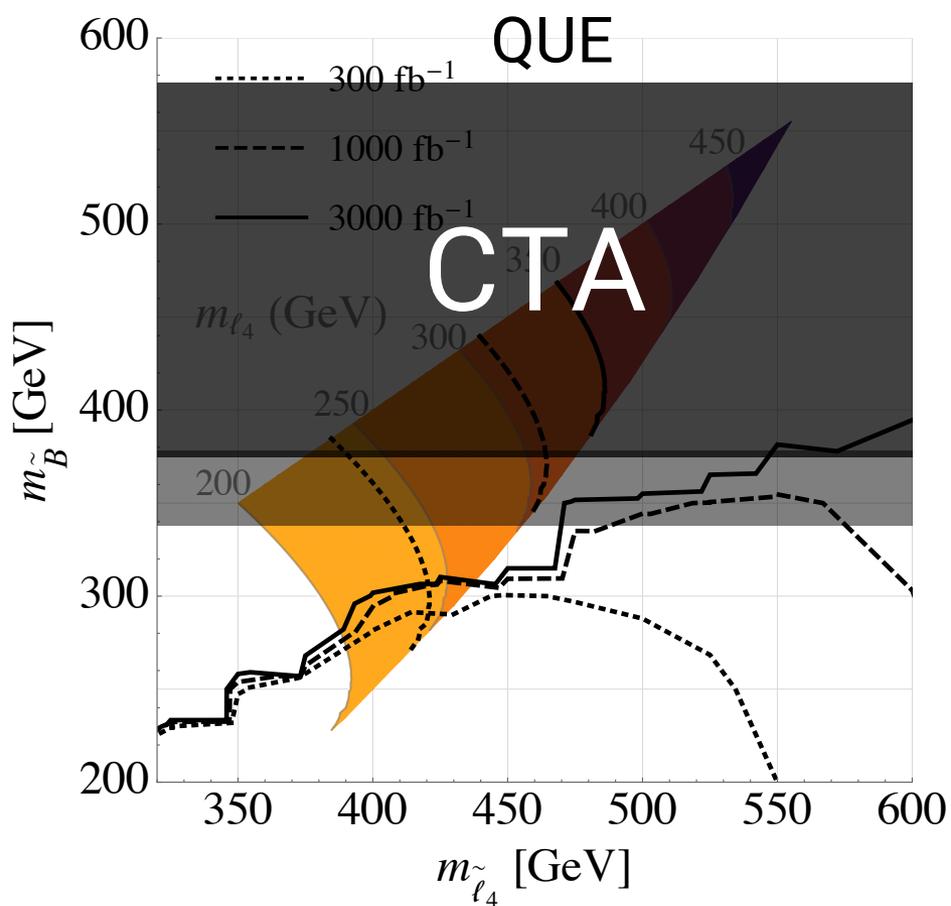
■ if 4G lepton decays to electron or muon



■ if 4G lepton decays to tau-lepton

LHC insensitive ... (· ω ·)

■ if 4G lepton decays to electron or muon



■ if 4G lepton decays to tau-lepton

LHC insensitive ... (· ω ·)

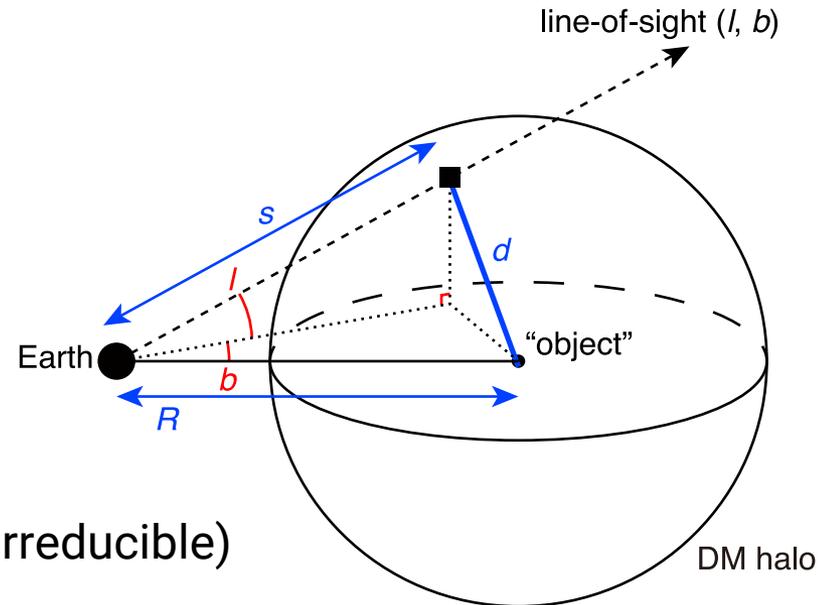
DM indirect detection Backup

■ charged particles → diffusion

- e : ~ 1 kpc are observable
- P : $\sim 0(10)$ kpc \sim Milky Way

■ neutral particles

- from (neighbor of) galactic center
 - larger density, huge BKG (miss-ID & irreducible)
 - $J \sim 10^{22}$ GeV²/cm⁵ (NFW; cuspy)
- from dwarf spheroidals (mini-galaxies near MW)
 - DM rich, less baryon → low BKG
 - $J < 10^{19-20}$ GeV²/cm⁵ (smaller profile dependence)



$$J = \int d\Omega_{l,b} \int_0^\infty ds \rho(d)^2$$

$$(d^2 = s^2 + R^2 - 2Rs \cos b \cos l)$$

■ < 100 GeV : satellites

- full-sky, $\sim 1\text{m}^2$, 5–10% energy resolution
- Fermi-LAT (2008) : gamma-ray to electron conversion

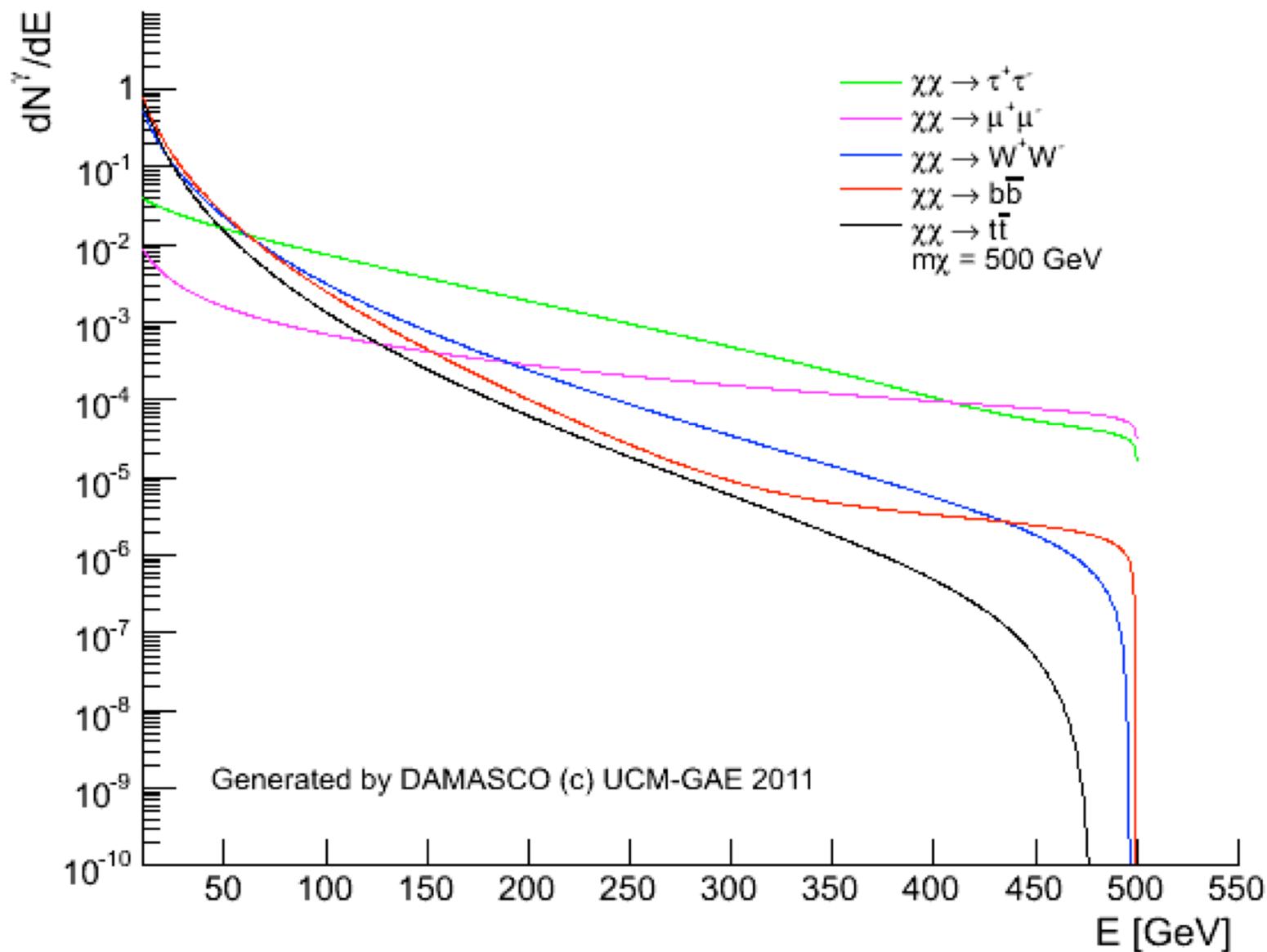
■ > 100 GeV : ground-based Air Cherenkov Telescopes

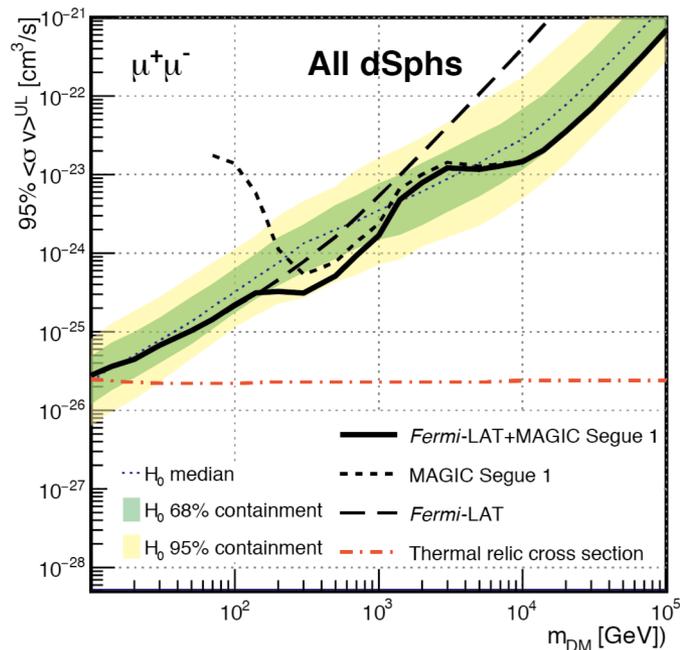
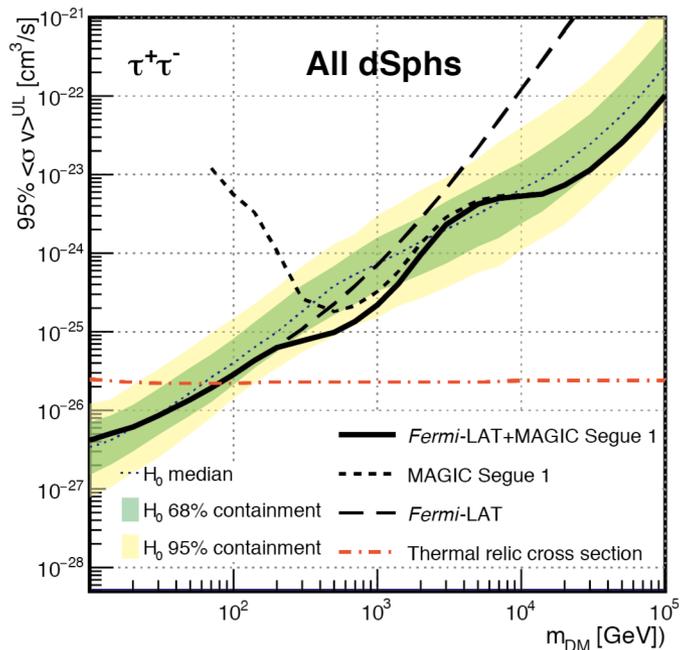
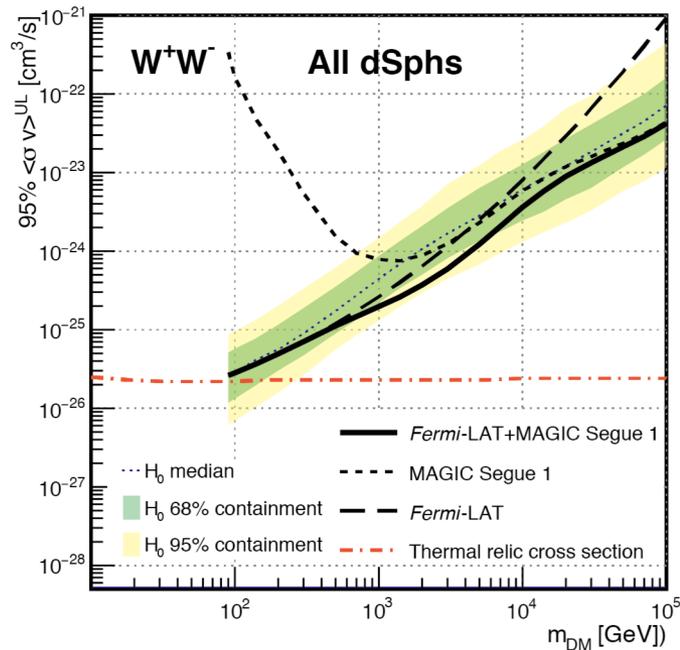
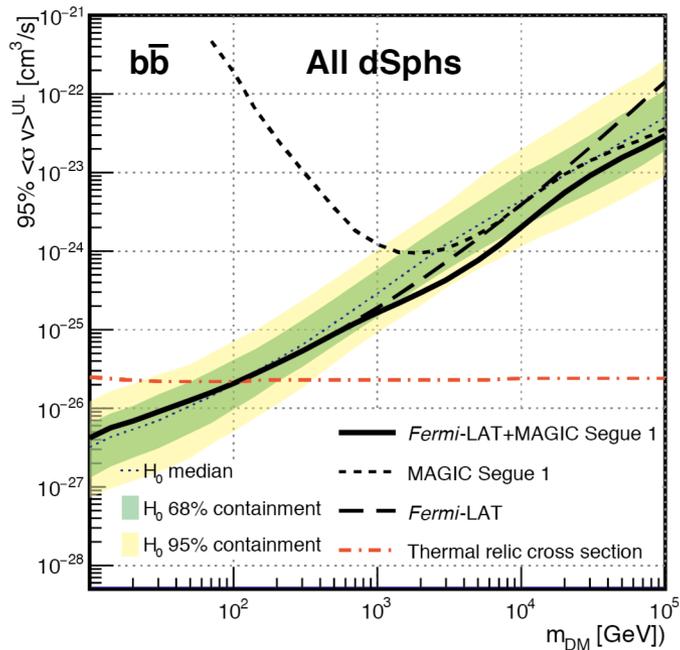
- several degree, 10^{5-6}m^2 , $\sim 20\%$ energy resolution
- VERITAS : 4x12m telescopes, Crab $36\sigma/\sqrt{\text{hr}} = 1\%\text{Crab}$ in 35h
- MAGIC : 2x17m telescope, $19\sigma/\sqrt{\text{hr}} = 2.2\%\text{Crab}$ in 50h
- HESS : 4x12m + 28m telescopes, $43\sigma/\sqrt{\text{hr}}$

■ > 10 TeV : ground-based Water Cherenkov

- HAWC : 2/3-sky, effective area similar to ACT but worse resolution

Spectra from Cembranos et al. (PRD 83:083507)

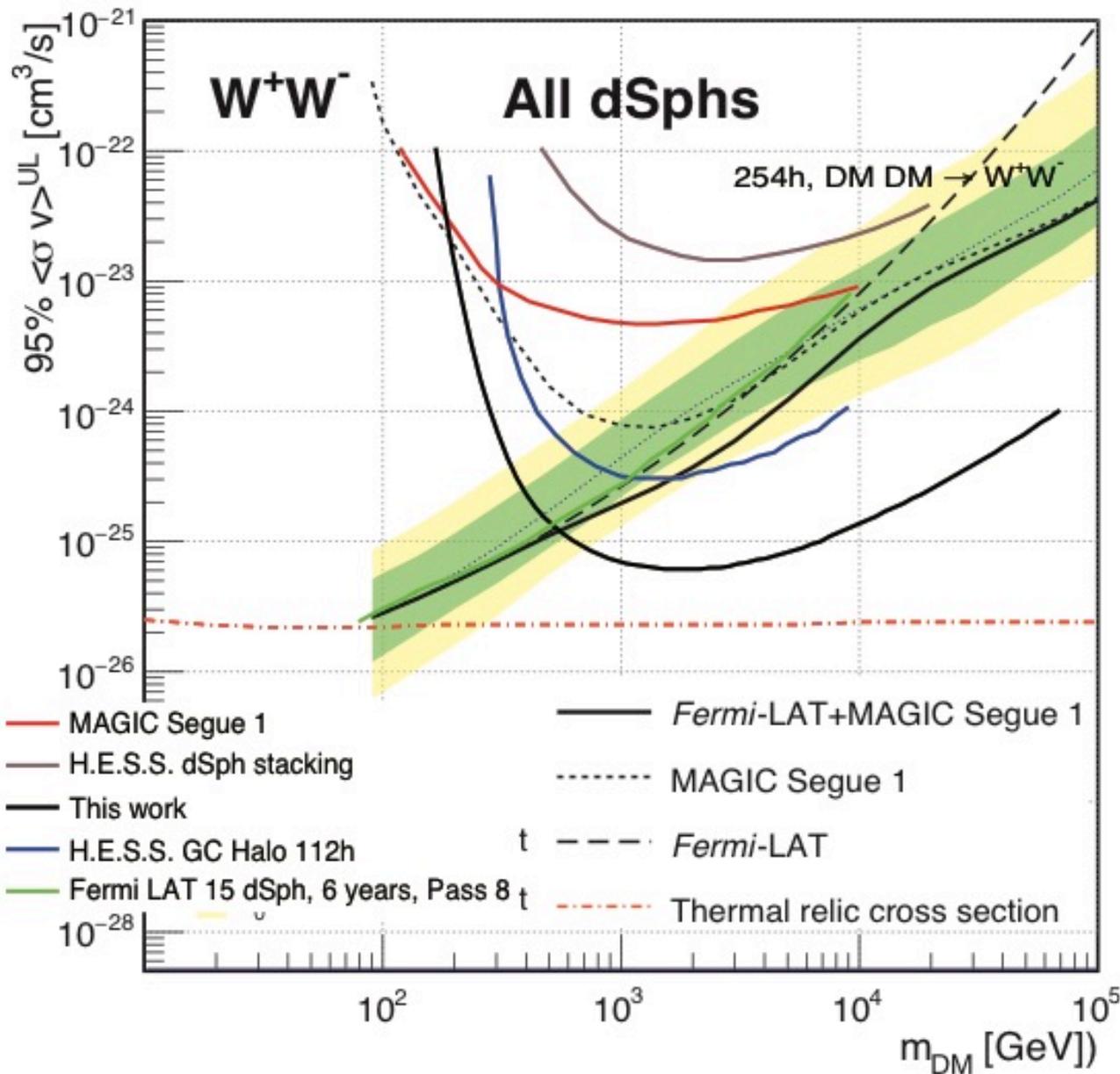
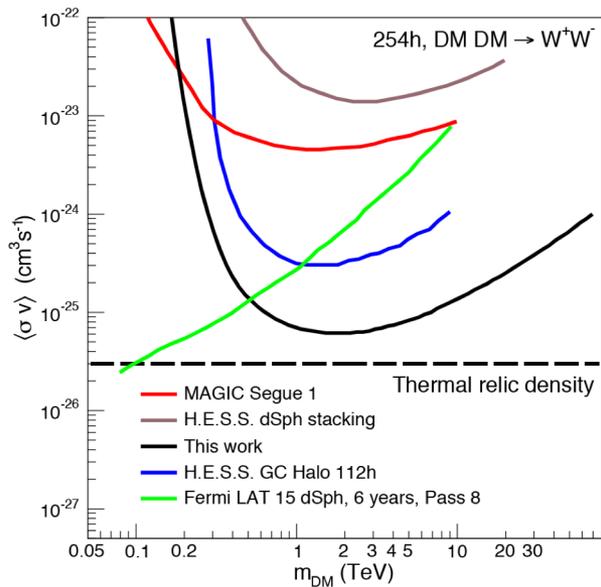
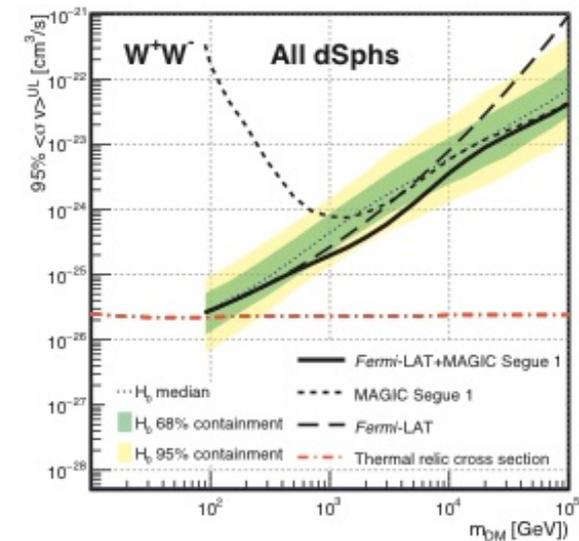




MAGIC:
158 hr of Segue 1

Fermi-LAT:
6 yr of 15 dSph
(incl. Segue 1)

DM profile: NFW



HESS assumes Einasto profile; for NFW weaker by factor ~2.

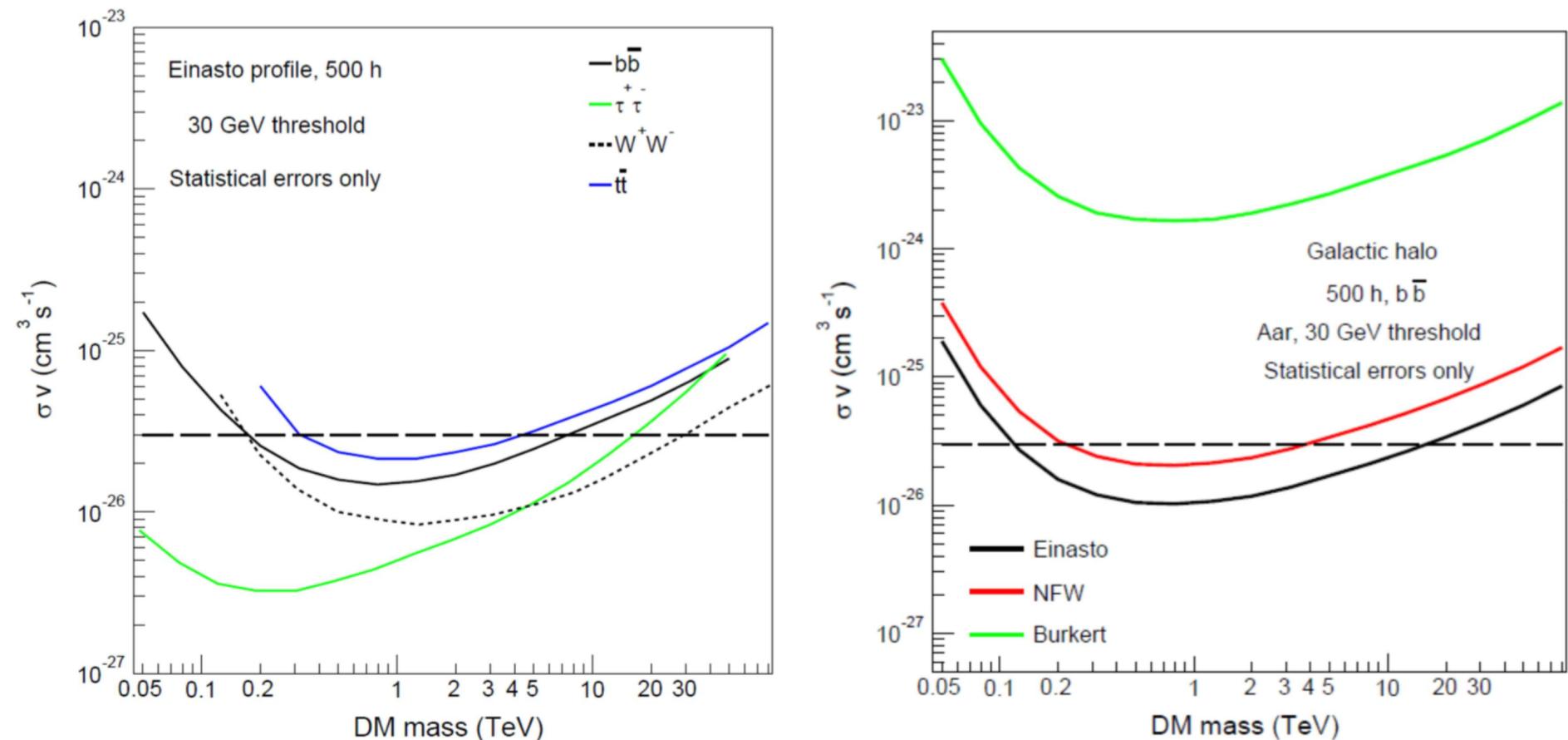
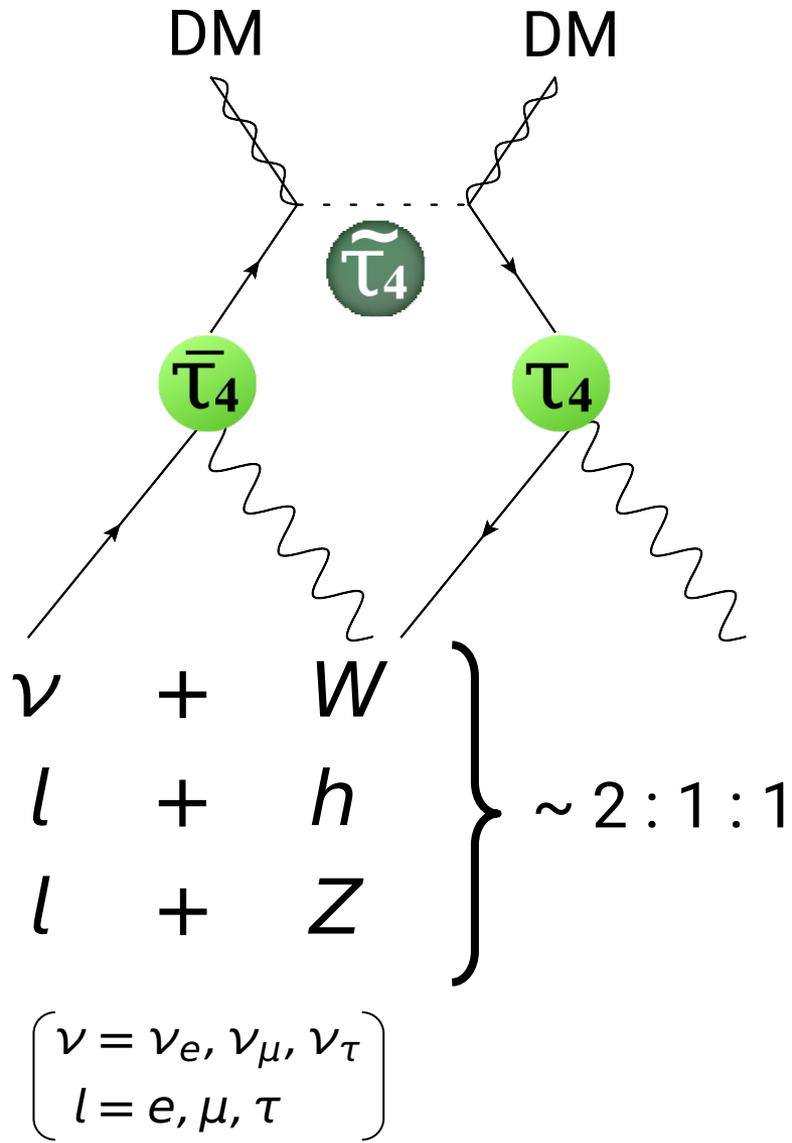
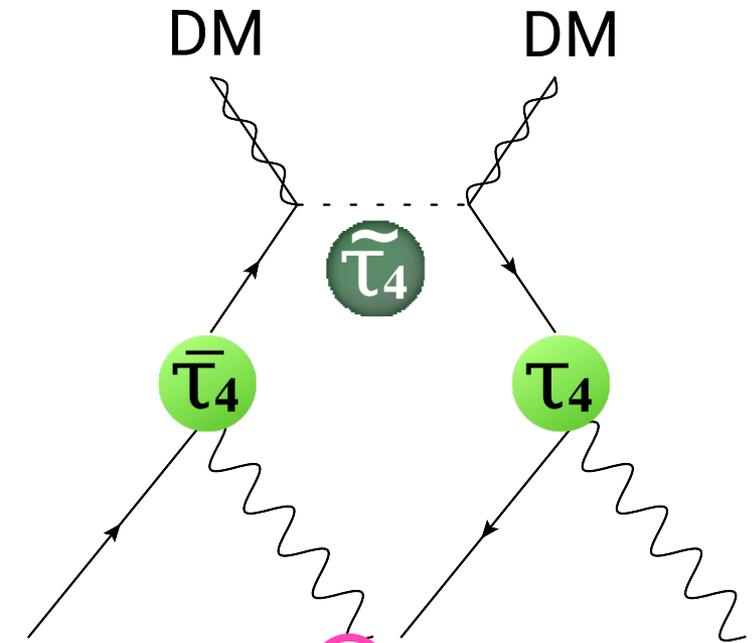


Figure 1. Left: Sensitivity for σv from observation on the Galactic Halo with Einasto dark matter profile and for different annihilation modes as indicated. **Right:** for cuspy (NFW, Einasto) and cored (Burkert) dark matter halo profiles. For both plots only statistical errors are taken into account. The dashed horizontal lines indicate the level of the thermal cross-section of $3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$.

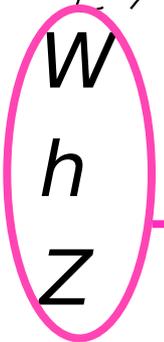
■ DM indirect detection by Gamma-ray observation



■ DM indirect detection by Gamma-ray observation



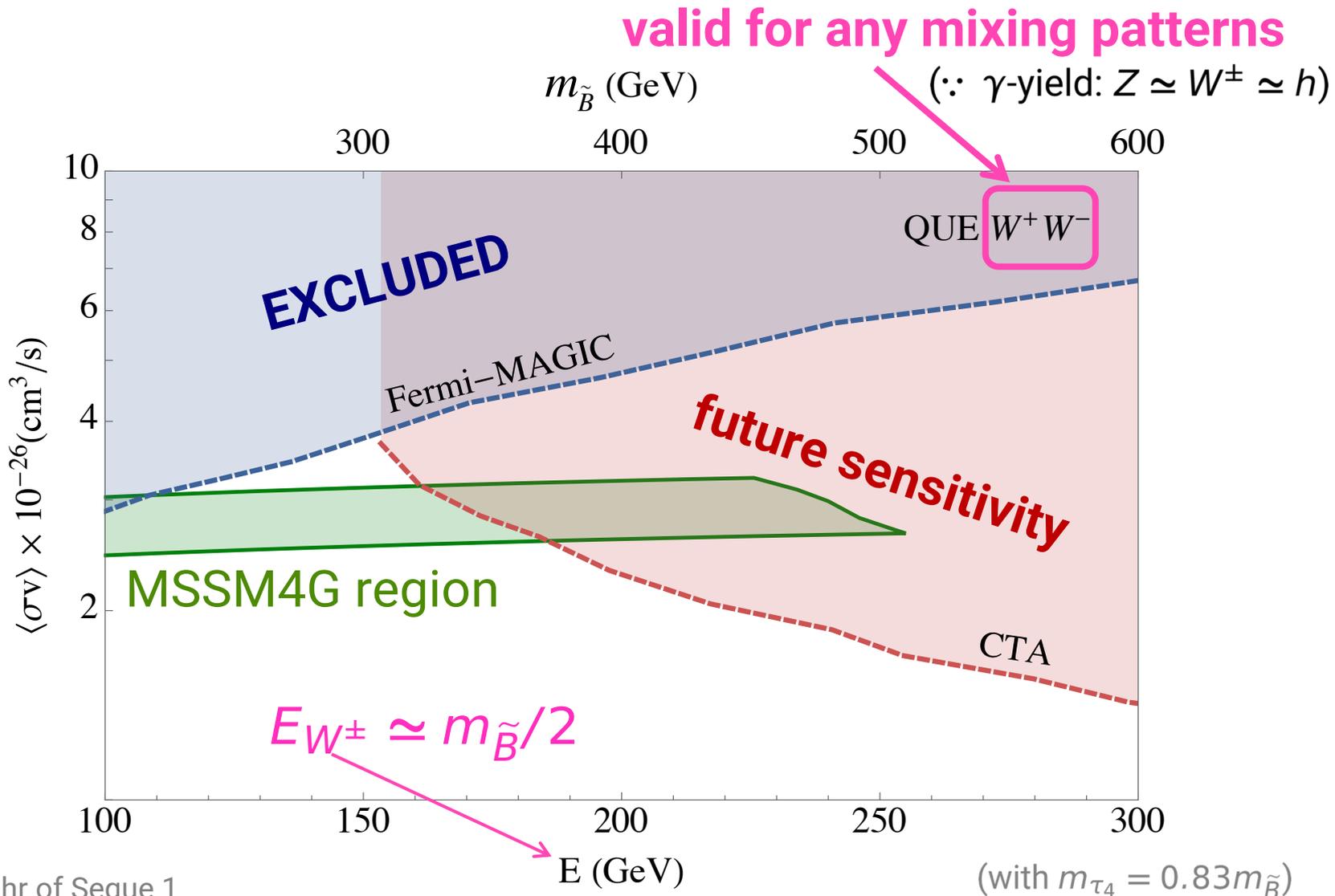
ν + W
 l + h
 l + Z



$\left(\begin{array}{l} \nu = \nu_e, \nu_\mu, \nu_\tau \\ l = e, \mu, \tau \end{array} \right)$



- Fermi-LAT (satellite)
- MAGIC (Air Cherenkov telescope)
- CTA (future A. C. Telescope)



MAGIC: 158 hr of Segue 1

Fermi-LAT: 6 yr of 15 dSph (incl. Segue 1)

DM profile: NFW

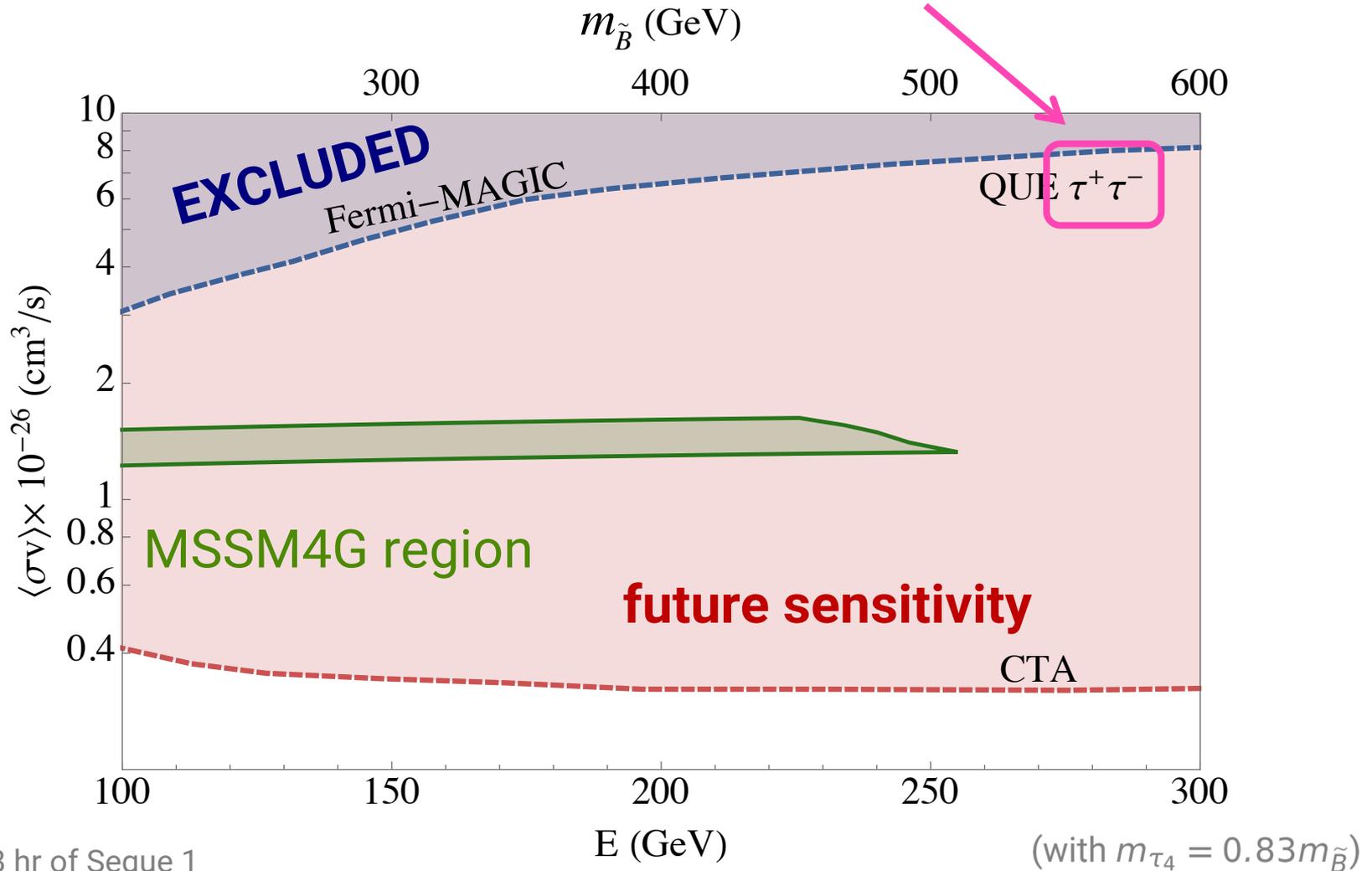
Fermi-LAT dominates MAGIC in almost all E -range.

CTA prospect : 500hr of Milky Way

DM profile: Einasto

No syst. unc. (stat only)

if 4G lepton decays to tau-lepton



MAGIC: 158 hr of Segue 1

Fermi-LAT: 6 yr of 15 dSph (incl. Segue 1)

DM profile: NFW

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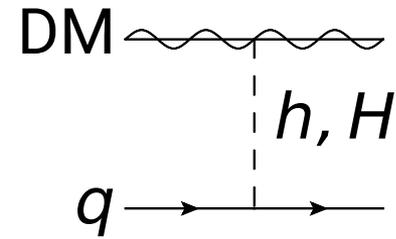
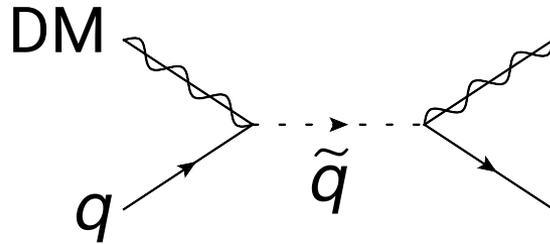
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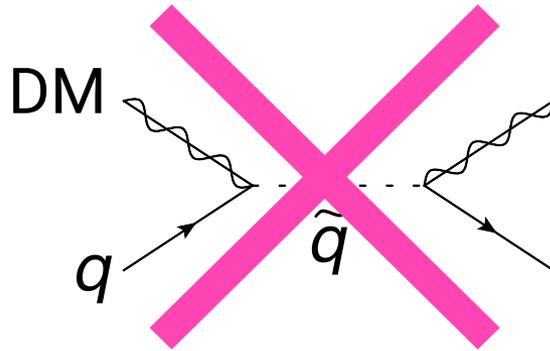
No syst. unc. (stat only)

SI direct detection constraints

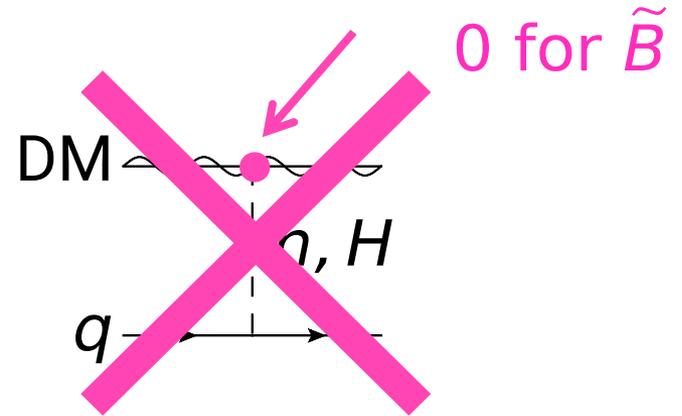
- spin-independent cross section $\sigma_{\text{SI}}(N\tilde{\chi}_1^0 \rightarrow N\tilde{\chi}_1^0)$



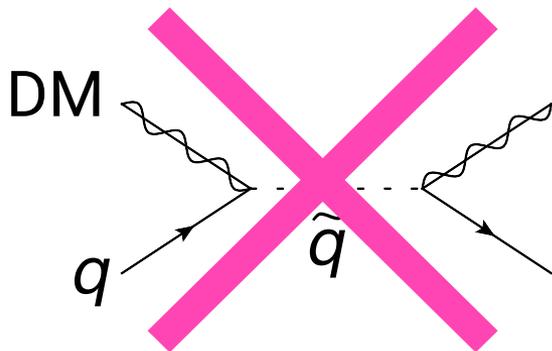
- spin-independent cross section $\sigma_{\text{SI}}(N\tilde{\chi}_1^0 \rightarrow N\tilde{\chi}_1^0)$



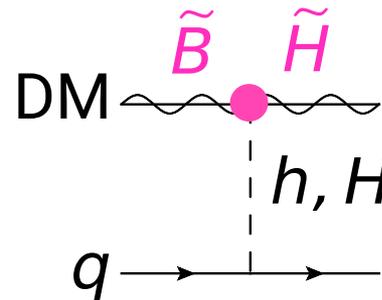
if \tilde{q} decoupled (as we assume)



- spin-independent cross section $\sigma_{\text{SI}}(N\tilde{\chi}_1^0 \rightarrow N\tilde{\chi}_1^0)$



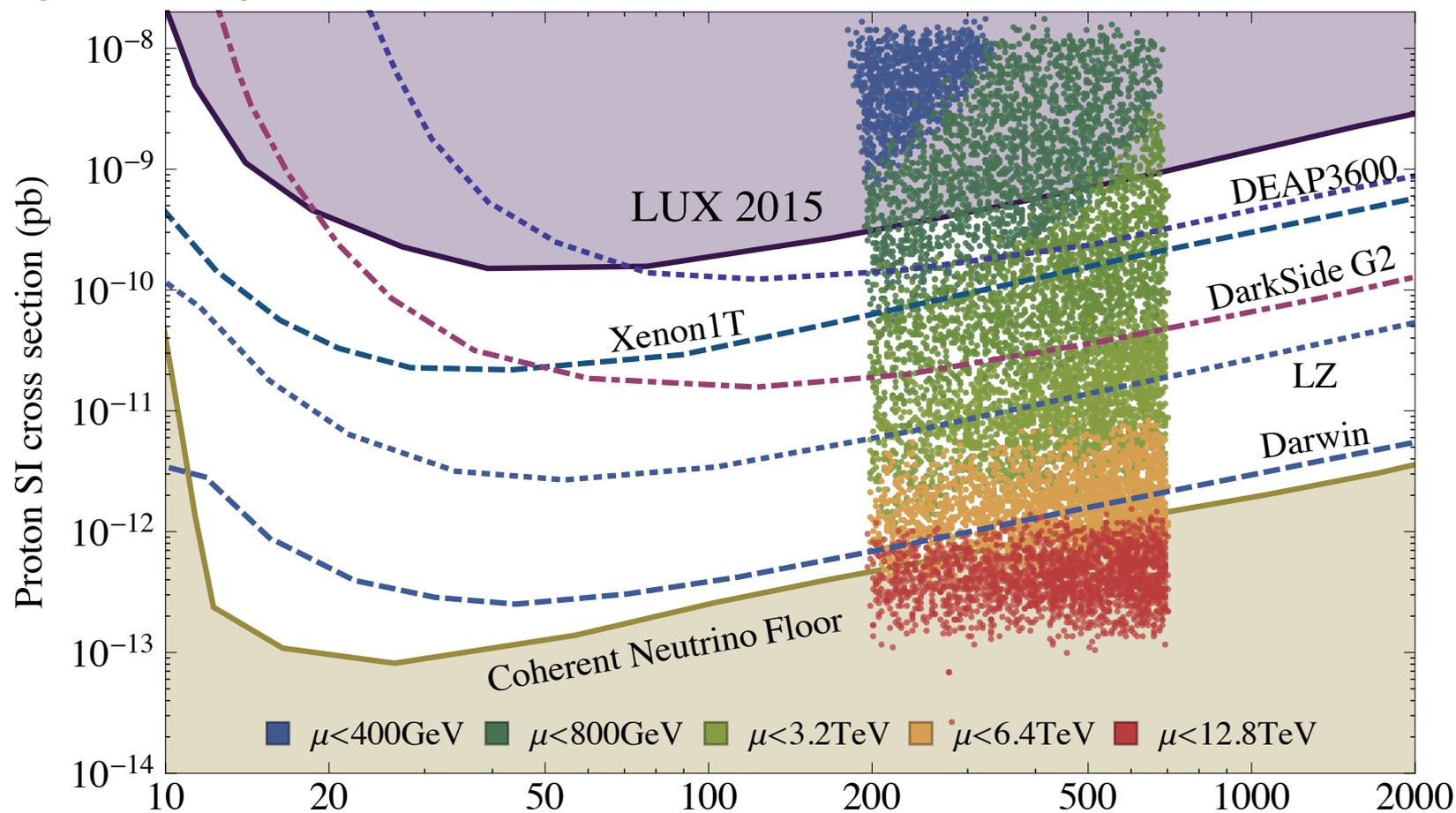
if \tilde{q} decoupled (as we assume)



$$\begin{aligned}
 \sigma_{\text{SI}} &\propto \left[N_{41} \left(N_{21} - N_{11} \tan \theta_w \right) \frac{g^2}{4m_W m_h^2} \right]^2 \\
 &\approx \left[\frac{M_1 m_Z s_w t_w}{\mu^2 - M_1^2 + m_Z^2 s_w^2} \frac{g^2}{4m_W m_h^2} \right]^2 \quad (M_1 \ll \mu \ll M_2, m_H)
 \end{aligned}$$

\tilde{W} in DM (~ 0) \tilde{B} in DM (~ 1)
 \tilde{H} in DM

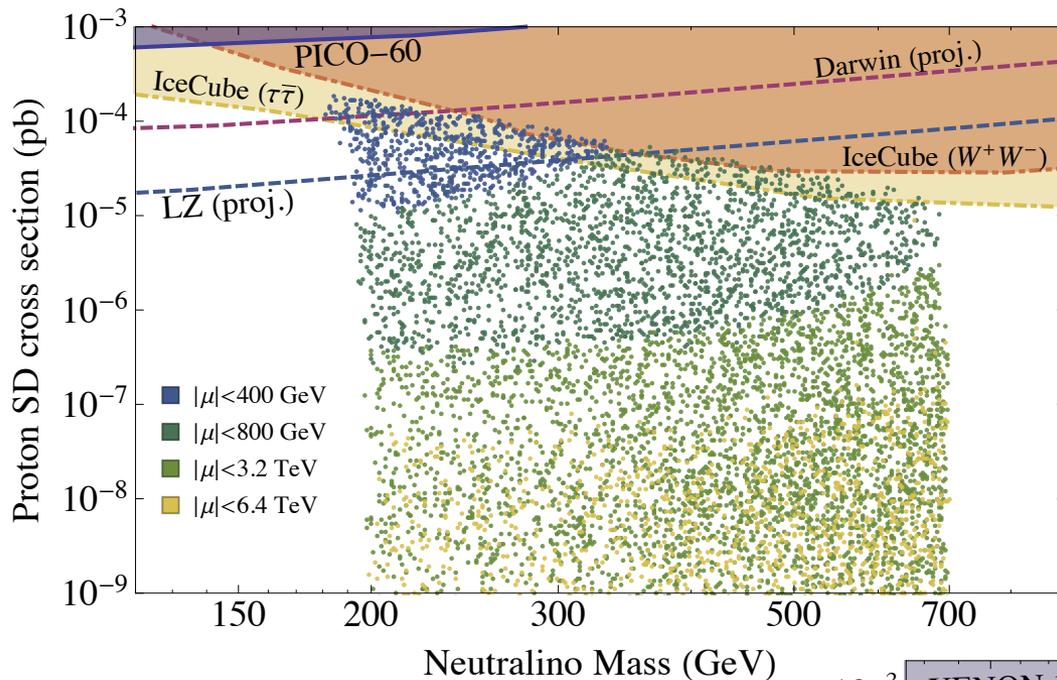
■ spin-independent cross section



$$\approx \left[\frac{M_1 m_Z s_w t_w}{\mu^2 - M_1^2 + m_Z^2 s_w^2} \frac{g^2}{4m_W m_h^2} \right]^2 \quad (M_1 \ll \mu \ll M_2, m_H)$$

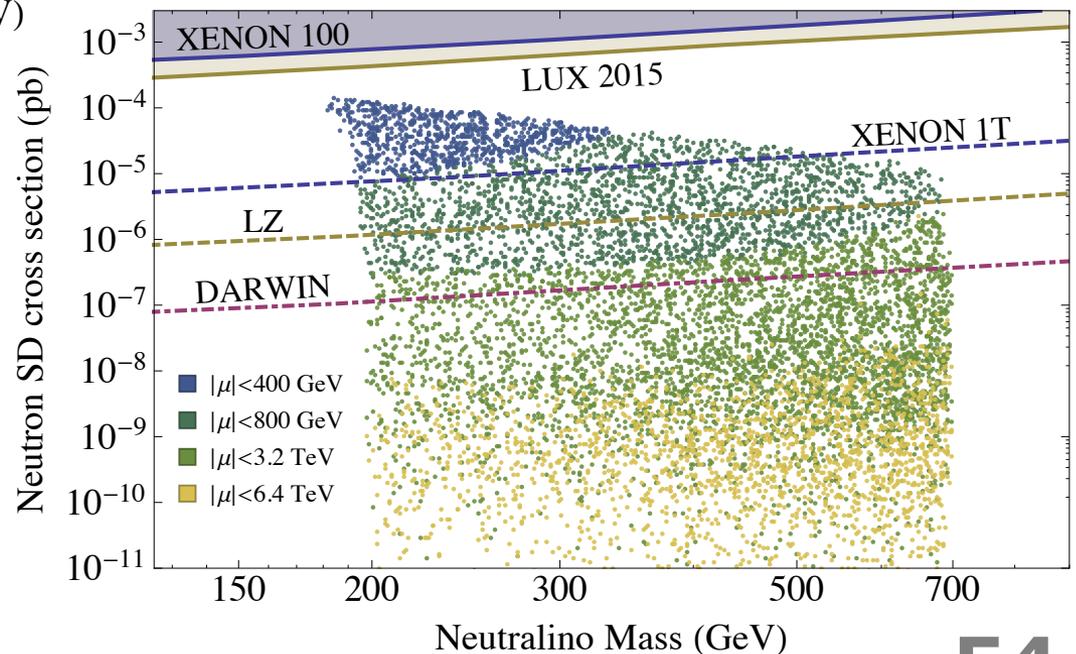
SD direct detection constraints

Spin-dependent direct detection constraints on MSSM4G scenario



Note

DM SI/SD constraints are the same as MSSM, but the parameter space is not preferred in the MSSM because of overabundance.



Vectorlike lepton search DATA

SR definition

	$WZ(j)$	$WZ(\ell)$	$ZZ(j)$	$ZZ(\ell)$
N_ℓ	≥ 3	≥ 4	≥ 4	≥ 5
N_j	≥ 2	< 2	≥ 2	—
$ m_{jj} - m_W $	$< 20 \text{ GeV}$	—	—	—
$ m_{jj} - m_Z $	—	—	$< 40 \text{ GeV}$	—
\cancel{E}_T	$> 60 \text{ GeV}$	$> 100 \text{ GeV}$	—	—
$N_{Z(\ell\ell)}$	—	—	≥ 1	≥ 1

TABLE IV: Selection flow of the background events in the vector-like lepton search. Upper bounds on the number of events in each SR, N_{UL} , are shown for three values of integrated luminosity, where systematic uncertainty of 20% as well as statistical uncertainty is included.

	background cross section [fb]				N_{UL}		
	di-boson	tri-boson	top	total	300 fb ⁻¹	1000 fb ⁻¹	3000 fb ⁻¹
$N_\ell \geq 3$	222	5.1	13.4	249	—	—	—
$WZ(j)^-$	0.071	0.013	0.082	0.166	25.1	70.4	200
$WZ(j)^Z$	0.643	0.071	0.183	0.898	111	359	1060
$WZ(\ell)^-$	0.014	0.025	0.017	0.056	11.9	27.4	71.1
$WZ(\ell)^Z$	< 0.001	0.005	0.003	0.008	5.1	7.9	14.5
$ZZ(j)^0$	0.194	0.016	0.058	0.268	37.2	111	321
$ZZ(j)^J$	0.064	0.007	0.022	0.093	16.4	41.8	114
$ZZ(j)^L$	0.182	0.012	0.024	0.218	31.2	91.7	263
$ZZ(j)^Z$	0.020	0.004	0.019	0.043	10.2	22.2	55.7
$ZZ(j)^{JL}$	0.060	0.005	0.009	0.075	14.2	35.3	94.3
$ZZ(j)^{JZ}$	0.008	0.001	0.008	0.017	6.7	11.9	25.6
$ZZ(j)^{LZ}$	0.020	0.004	0.019	0.043	10.2	22.2	55.9
$ZZ(j)^{JLZ}$	0.008	0.001	0.008	0.017	6.7	11.9	25.5
$ZZ(\ell)$	< 0.001	0.005	< 0.001	0.005	4.7	6.8	11.5
$ZZ(\ell)^{<2}$	< 0.001	0.003	< 0.001	0.004	4.2	5.8	9.2
$ZZ(\ell)^{<1}$	< 0.001	0.001	< 0.001	0.001	3.6	4.5	6.3

Z-flag for $WZ(j)$: a Z-like $\ell\ell$ (SFOS, $|m_{\ell\ell} - m_Z| < 10$ GeV)

Z-flag for $WZ(\ell)$: a Z-like $\ell\ell$ in 3rd-leading leptons

J-flag for $ZZ(j)$: a Z-like jj (10 GeV)

L-flag for $ZZ(j)$: a Z-like $\ell\ell$ in 2nd-leading leptons

Z-flag for $ZZ(j)$: leading-lepton does NOT form Z-like pairs

$ZZ(\ell)$ divided by number of jets

TABLE V: Selection flow of the signal events in searches for the e - or μ -mixed τ_4 in the QUE model, displayed as a signal cross section in fb. SRs marked with *, † and ‡ are the most sensitive for exclusion at $\mathcal{L} = 300, 1000, \text{ and } 3000 \text{ fb}^{-1}$, respectively.

m_τ [GeV], mixing	200, e	200, μ	300, e	300, μ	400, e	400, μ
total	95.7	96.0	21.2	21.2	6.76	6.74
$N_\ell \geq 3$	2.23	2.42	0.634	0.671	0.231	0.230
$WZ(j)^-$	0.018	0.022	0.020	0.024	0.011	0.012
$WZ(j)^Z$	0.049	0.063	0.034	0.036	0.014	0.014
$WZ(\ell)^Z$	0.012	0.014	0.008‡	0.008	0.003	0.004‡
$ZZ(j)^0$	0.066	0.065	0.035	0.044	0.015	0.015
$ZZ(j)^J$	0.035	0.033	0.018	0.023	0.008	0.007
$ZZ(j)^L$	0.045	0.048	0.026	0.031	0.011	0.012
$ZZ(j)^Z$	0.039*	0.042*	0.025*†	0.029†	0.010*	0.012†
$ZZ(j)^{JL}$	0.025	0.025	0.013	0.016	0.006	0.006
$ZZ(j)^{JZ}$	0.021	0.022	0.013	0.015‡	0.005	0.006
$ZZ(j)^{LZ}$	0.039	0.042	0.025	0.029*	0.010†	0.012*
$ZZ(j)^{JLZ}$	0.021	0.022	0.013	0.015	0.005	0.006
$ZZ(\ell)$	0.015†‡	0.014†‡	0.005	0.007	0.003‡	0.002
$ZZ(\ell)^{<2}$	0.010	0.009	0.003	0.004	0.002	0.001
$ZZ(\ell)^{<1}$	0.004	0.003	0.001	0.002	8×10^{-4}	6×10^{-4}

Z-flag for $WZ(j)$: a Z -like $\ell\ell$ (SFOS, $|m_{\ell\ell} - m_Z| < 10 \text{ GeV}$)

Z-flag for $WZ(\ell)$: a Z -like $\ell\ell$ in 3rd-leading leptons

J-flag for $ZZ(j)$: a Z -like jj (10 GeV)

L-flag for $ZZ(j)$: a Z -like $\ell\ell$ in 2nd-leading leptons

Z-flag for $ZZ(j)$: leading-lepton does NOT form Z -like pairs

$ZZ(\ell)$ divided by number of jets

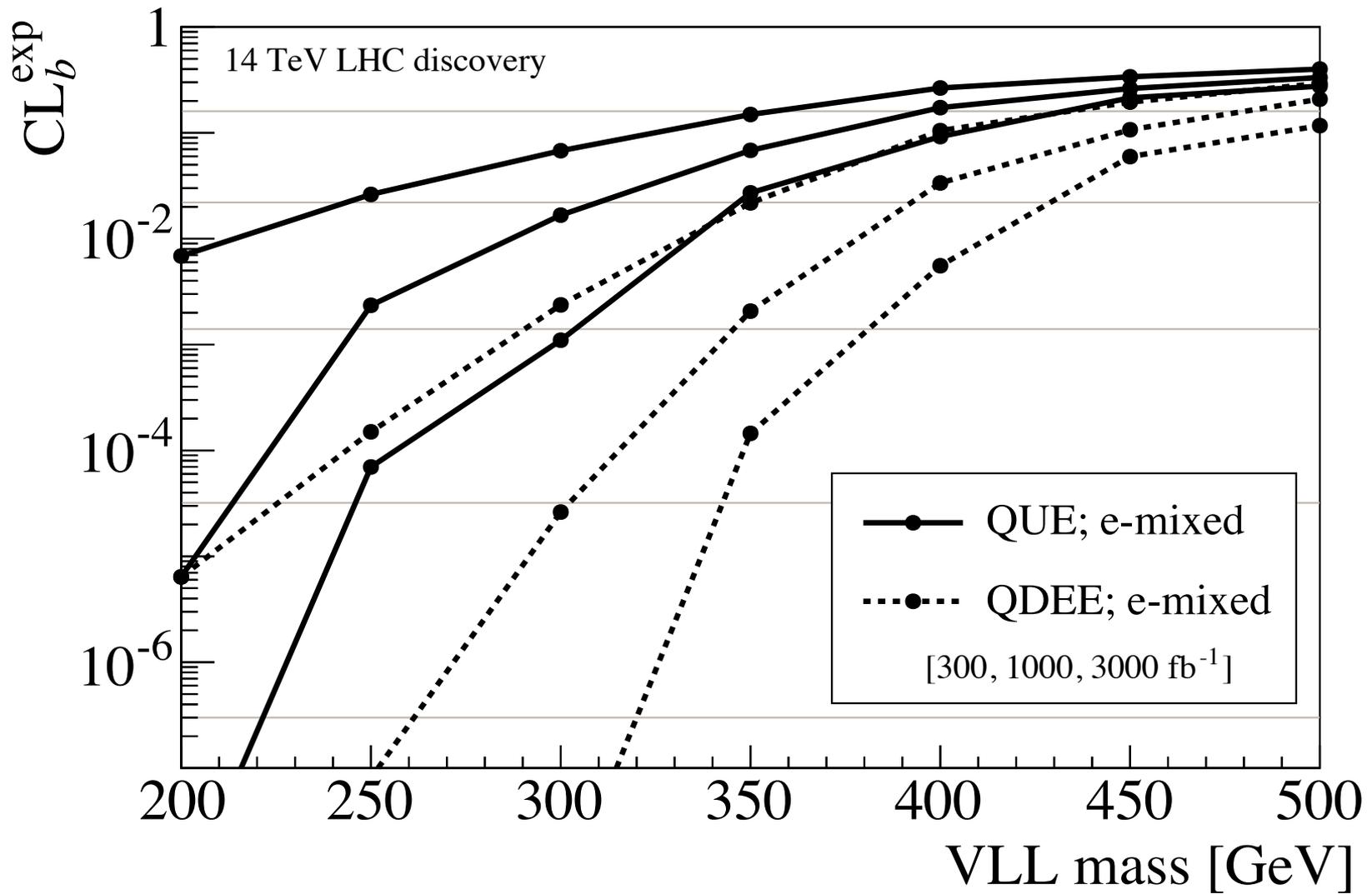
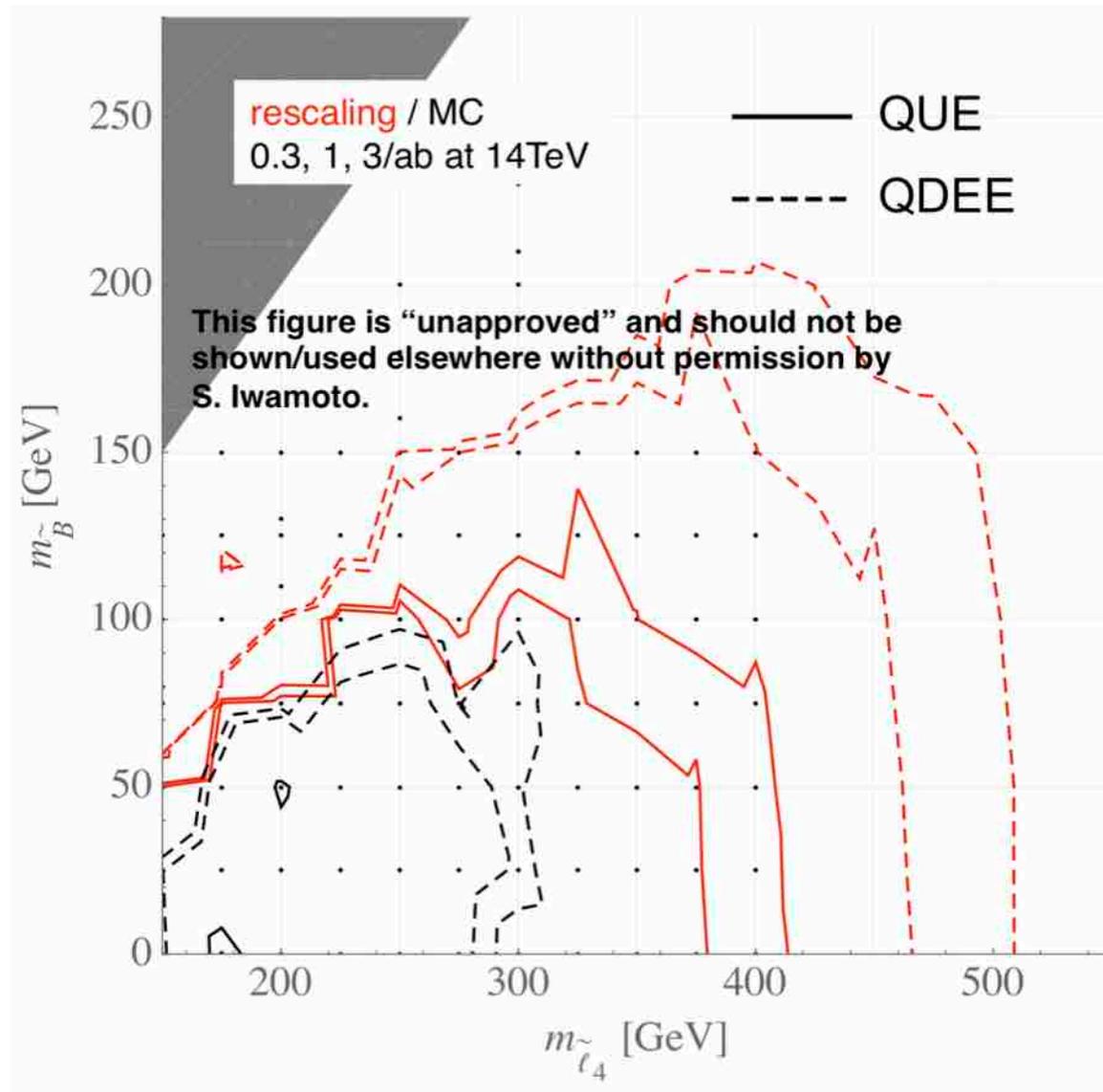


TABLE II: Future prospects for searches for vector-like leptons at the 14 TeV LHC for three values of integrated luminosity. The first table is for the QUE models, and the second for the QDEE models. We consider vector-like leptons with a mass $m_{\ell_4} \geq 200$ GeV; the expressions 0^{+250} GeV etc. show that the central value of exclusion or discovery limit is below our model points and we may achieve the limit of 250 GeV with 1σ statistical fluctuation. In the dashed entries the upper limit is less than 200 GeV even with 1σ statistical fluctuation. The CL_s method is used for statistical treatment, where the statistical uncertainty and a 20% systematic uncertainty for the background contribution are taken into account, while the theoretical uncertainty on the signal cross section as well as the NLO correction are not considered. See Appendix B for further details.

QUE model		300 fb ⁻¹	1000 fb ⁻¹	3000 fb ⁻¹
95% CL exclusion	<i>e</i> -mixed	240 ⁺⁶⁰ GeV	310 ⁺⁵⁰ ₋₆₀ GeV	350 ⁺⁴⁰ ₋₄₀ GeV
	μ -mixed	270 ⁺⁵⁰ GeV	330 ⁺⁴⁰ ₋₆₀ GeV	370 ⁺⁴⁰ ₋₄₀ GeV
3 σ discovery	<i>e</i> -mixed	0 ⁺²⁵⁰ GeV	250 ⁺⁶⁰ ₋₄₀ GeV	300 ⁺⁵⁰ ₋₅₀ GeV
	μ -mixed	0 ⁺²⁸⁰ GeV	260 ⁺⁷⁰ ₋₆₀ GeV	320 ⁺⁵⁰ ₋₄₀ GeV
5 σ discovery	<i>e</i> -mixed	—	0 ⁺²¹⁰ GeV	220 ⁺²⁰ ₋₂₀ GeV
	μ -mixed	—	0 ⁺²¹⁰ GeV	240 ⁺²⁰ ₋₂₀ GeV

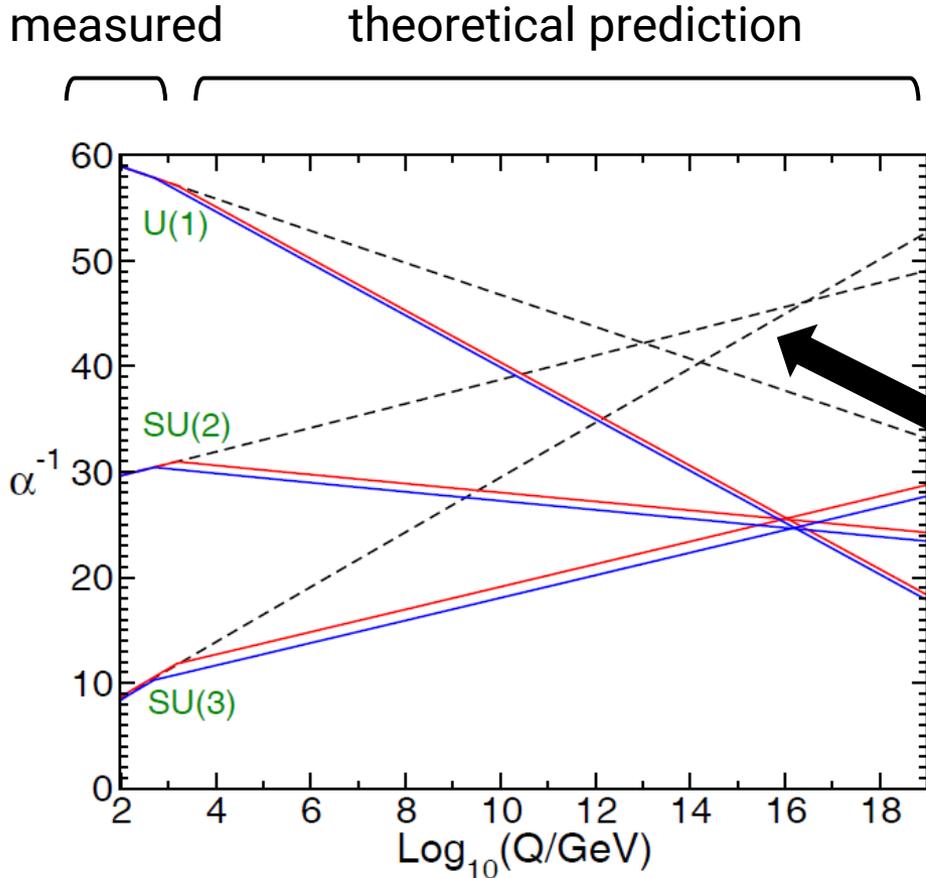
QDEE model		300 fb ⁻¹	1000 fb ⁻¹	3000 fb ⁻¹
95% CL exclusion	<i>e</i> -mixed	350 ⁺⁴⁰ ₋₅₀ GeV	390 ⁺⁴⁰ ₋₄₀ GeV	430 ⁺⁴⁰ ₋₄₀ GeV
	μ -mixed	360 ⁺⁴⁰ ₋₄₀ GeV	400 ⁺⁴⁰ ₋₄₀ GeV	440 ⁺⁴⁰ ₋₄₀ GeV
3 σ discovery	<i>e</i> -mixed	290 ⁺⁶⁰ ₋₇₀ GeV	340 ⁺⁶⁰ ₋₄₀ GeV	380 ⁺⁵⁰ ₋₄₀ GeV
	μ -mixed	310 ⁺⁶⁰ ₋₅₀ GeV	360 ⁺⁴⁰ ₋₃₀ GeV	400 ⁺⁴⁰ ₋₃₀ GeV
5 σ discovery	<i>e</i> -mixed	0 ⁺²⁰⁰ GeV	260 ⁺⁴⁰ ₋₅₀ GeV	310 ⁺²⁰ ₋₃₀ GeV
	μ -mixed	0 ⁺²⁶⁰ GeV	280 ⁺³⁰ ₋₃₀ GeV	320 ⁺⁴⁰ ₋₂₀ GeV



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Gauge coupling unification

- SM \ni 3 forces : U(1), SU(2), SU(3) [Why three?]



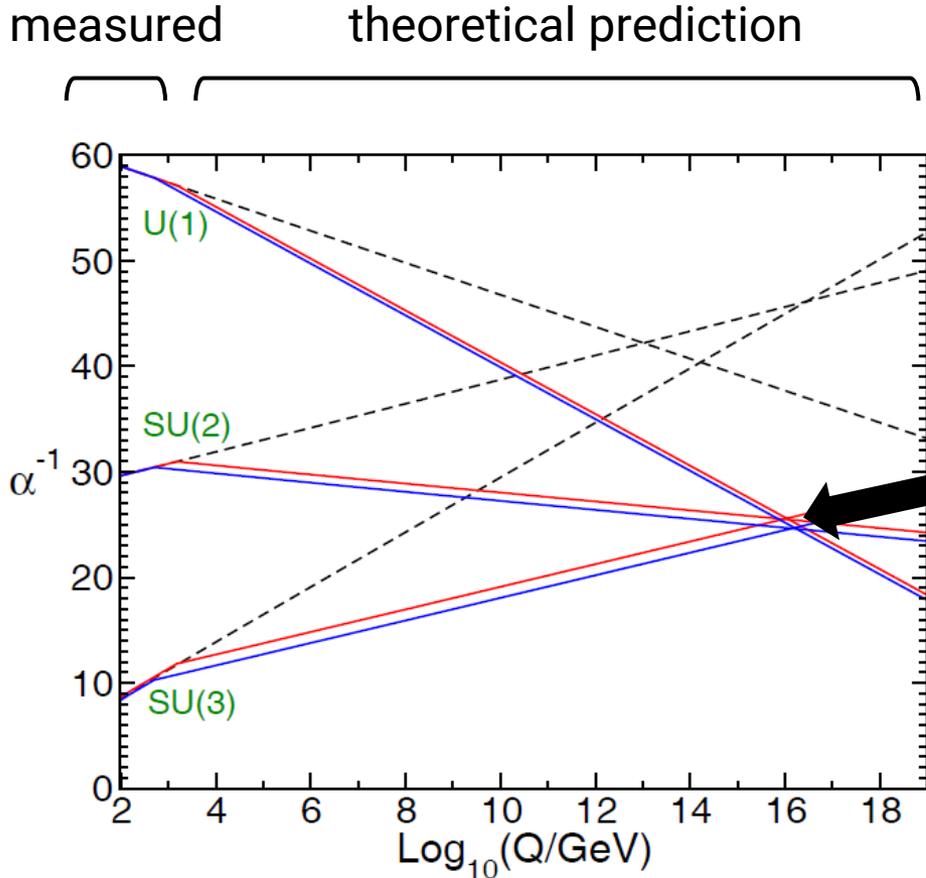
Gauge coupling unification

	mass \rightarrow $\approx 2.3 \text{ MeV}/c^2$	$\approx 1.275 \text{ GeV}/c^2$	$\approx 173.07 \text{ GeV}/c^2$	0	$\approx 126 \text{ GeV}/c^2$
charge \rightarrow 2/3	u	c	t	g	H
spin \rightarrow 1/2	up	charm	top	gluon	Higgs boson
	$\approx 4.8 \text{ MeV}/c^2$	$\approx 95 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	0	0
-1/3	d	s	b	γ	0
1/2	down	strange	bottom	photon	0
	$0.511 \text{ MeV}/c^2$	$105.7 \text{ MeV}/c^2$	$1.777 \text{ GeV}/c^2$	0	$91.2 \text{ GeV}/c^2$
-1	e	μ	τ	Z	0
1/2	electron	muon	tau	Z boson	1
	$< 2.2 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 15.5 \text{ MeV}/c^2$	0	$80.4 \text{ GeV}/c^2$
0	ν_e	ν_μ	ν_τ	W	± 1
1/2	electron neutrino	muon neutrino	tau neutrino	W boson	1

Figure from S. P. Martin, *A Supersymmetry Primer*, hep-ph/9709356

Gauge coupling unification

- SM \ni 3 forces : U(1), SU(2), SU(3) [Why three?]



Gauge coupling unification

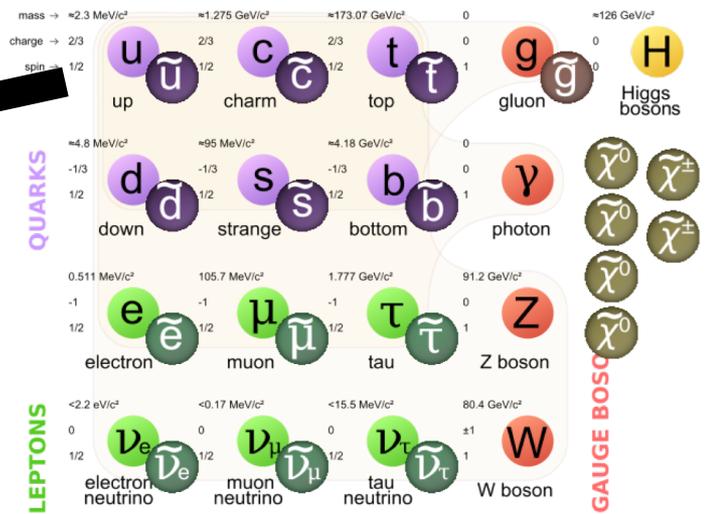


Figure from S. P. Martin, *A Supersymmetry Primer*, hep-ph/9709356